HAUSDORFF DIMENSION OF LIMIT SETS FOR PARABOLIC IFS WITH OVERLAPS

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ABSTRACT. We study parabolic iterated function systems with overlaps on the real line. We show that if a d-parameter family of such systems satisfies a transversality condition, then for almost every parameter value the Hausdorff dimension of the limit set is the minimum of 1 and the least zero of the pressure function. Moreover, the local dimension of the exceptional set of parameters is estimated. If the least zero is greater than 1, then the limit set (typically) has positive Lebesgue measure. These results are applied to some specific families including those arising from a class of continued fractions.

1. INTRODUCTION

Let $\Phi = \{\phi_1, \ldots, \phi_k\}$ be a collection of self-maps on a closed interval $X \subset \mathbb{R}$. We call Φ an **iterated function system** (IFS). Under standard contractivity hypotheses, there is a unique non-empty compact set J_Φ such that $J_\Phi = \bigcup_{j=1}^k \phi_j(J_\Phi)$, called the **limit set**, or attractor, of the IFS.

If the sets $\phi_j(X)$ are mutually disjoint, then J_Φ is a Cantor set. If, in addition, ϕ_j are monotone, the limit set is known as a "cookie-cutter"; then it is more common to view J_Φ as the repeller of an expanding map $f : \bigcup_{j=1}^k \phi_j(X) \to X$ defined by $f(x) = \phi_j^{-1}(x)$ for $x \in \phi_j(X)$. Suppose that all the maps ϕ_j are in $\mathcal{C}^{1+\theta}(X)$ for some $\theta \in (0, 1]$ and are **hyperbolic**, that is, $0 < |\phi_j'(x)| < 1$ on X. Then the Hausdorff dimension $\dim_{\mathrm{H}}(J_\Phi)$ is given by Bowen's formula [Bo2, R]:

$$\dim_{\mathrm{H}}(J_\Phi) = s(\Phi) \qquad \text{where} \qquad P_\Phi(s(\Phi)) = 0.$$

Here $P_\Phi(t)$ is the pressure function, which can be defined by

$$P_\Phi(t) = \lim_{n\to\infty} \frac{1}{n} \log \sum_{\omega\in I^n} \| \phi_\omega' \|^t$$

where $I = \{1, \ldots, k\}$, $\phi_\omega = \phi_{\omega_1} \circ \cdots \circ \phi_{\omega_n}$, and $\| \cdot \|$ is the supremum norm on X.

The dimension formula was extended in [U] to the **parabolic** case, where some of the maps have a neutral fixed point. A parabolic IFS is not uniformly contracting, which

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makes the situation more subtle. It was proved in [U] that the Hausdorff dimension of the limit set is the least zero of the pressure function. In contrast with the hyperbolic case (where the zero is unique), in the parabolic case the pressure function is identically zero for all t larger than the least zero.

We study the case when the "pieces" of the limit set are allowed to overlap. One-dimensional IFS with overlaps arise naturally in the study of higher-dimensional dynamical systems [PrU, BU, Si1, Si2, SSo] (in this sentence the "dimension" refers to the phase space rather than to the limit set). Moreover, IFS with overlaps occur in some problems on random matrix products, random continued fractions, and in prediction theory, see [Pi, Ly, LL].

There are many open problems related to IFS with overlaps, which are notoriously difficult even when the maps are linear, see [PSo3]. The dimension of the limit set may be strictly less than the least zero of the pressure function, for instance, if $\phi_\omega \equiv \phi_\tau$ for two distinct words ω and τ. Since it is often hard to analyze an individual system, one can try to investigate a "typical" (in the sense of Lebesgue measure) IFS in a parameterized family. This method was first used by Falconer [F1] who considered families of linear contractions with a linear dependence on parameter. Further work in this direction was done in [PoS, So1, PSo2, So2, SSo]. An important role in these papers was played by a certain "transversality condition", which controls the way the IFS depends on parameters.

In this paper this approach is extended to a class of non-linear IFS. Our main result (Theorem 6.1) states that if a parameterized family $\{\Phi^{\mathbf{t}}\}$ of parabolic IFS satisfies the transversality condition, then for Lebesgue-a.e. parameter $\mathbf{t}$ the Hausdorff dimension of the limit set is given by

$$\dim_{\mathrm{H}}(J_{\Phi^{\mathbf{t}}}) = \min\{1, s(\mathbf{t})\} \quad \text{where} \quad s(\mathbf{t}) := \min\{s : \ P_{\Phi^{\mathbf{t}}}(s) = 0\}.$$

Moreover, the limit set $J_{\Phi^{\mathbf{t}}}$ has positive Lebesgue measure for a.e. $\mathbf{t}$ such that $s(\mathbf{t}) > 1$. If a slightly stronger version of transversality is imposed, then the local dimension of the exceptional set in the first statement can be estimated above by $s(\mathbf{t}) + (d-1)$, where d is the number of independent parameters involved.

We illustrate our results by the following example (see Corollaries 7.4 and 7.5). Let $\phi(x) = \frac{x}{1+x}$. This function is parabolic on $[0,1]$. Let $A = \{a_1, \ldots, a_k\}$, with $k \geq 3$, and consider the family of IFS on $[0,1]$

$$\Phi_A = \{\phi(x + a_j)\}_{j=1}^k \quad \text{for} \quad A \in U := \{A \in \mathbb{R}^k : \ a_k = 0, \ a_j > 0, \ j = 1, \ldots, k-1\}.$$

Denote by J_A the limit set of the IFS Φ_A and let $s(A) = \min\{t > 0 : \ P_{\Phi_A}(t) = 0\}$. We will prove that $s : \ U \to \mathbb{R}$ is a continuous function, so $U_{<1} = \{A \in U : \ s(A) < 1\}$ and $U_{>1} = \{A \in U : \ s(A) > 1\}$ are open sets.

Proposition 1.1. **(i)** *For Lebesgue-a.e.* $A \in U$,

$$\dim_{\mathrm{H}}(J_A) = \min\{s(A), 1\}.$$

(ii) *For any subset* $G \subset U_{<1}$ *we have*

$$\dim_{\mathrm{H}}\{A \in G : \ \dim_{\mathrm{H}}(J_A) < s(A)\} \le \sup_G s(A) + (k-2).$$

(iii) *For Lebesgue-a.e.* $A \in U_{>1}$ *the set* J_A *has positive Lebesgue measure.*

(iv) *Similar results hold for the one-parameter family* Φ_A *where* $A = \{a, 2, 0\}$ *and* $a \in (0, 2)$.

Remarks. 1. We assumed that $k \ge 3$ since for $k = 2$ either the limit set is an interval, or the IFS has no overlaps, so the result is true for all parameters by [U].

2. Proposition 1.1(iii) concerning the positive measure of the limit set reflects a phenomenon which cannot occur in the non-overlapping case. It is an open problem whether such limit sets can be "fat" Cantor sets or they necessarily contain intervals.

3. The limit set J_A can be described as the set of continued fractions of the form

$$y = [1, Y_1, 1, Y_2, 1, Y_3, \dots] = \cfrac{1}{1 + \cfrac{1}{Y_1 + \cfrac{1}{1 + \dots}}}$$

where $Y_i \in A$. The dimension of sets arising by some restriction in their expansions (continued fractions, λ-expansions, etc.) was studied by many authors. IFS with overlaps arise when the expansion for some numbers is non-unique. The family of linear IFS $\{\lambda x,\ \lambda x + 1,\ \lambda x + 3\}$, with $\lambda \in (\frac{1}{4}, \frac{2}{5})$, investigated in [KSS, PoS, So1], was an important "testing ground" in the study of IFS with overlaps. The family $\{\phi(x), \phi(x+\alpha), \phi(x+2)\}$, with $\alpha \in (0, 2)$, that we consider in Proposition 1.1(iv) is a non-linear parabolic analog.

4. A related problem is to analyze invariant (stationary) measures on the limit set of an IFS. The fundamental question is whether this measure is singular or absolutely continuous. This is interesting already for $k = 2$, when the limit set is an interval. R. Lyons [Ly] investigated a family of such measures for the IFS Φ_A with $A = \{\alpha, 0\}$. He showed singularity for a certain interval of parameters and asked if the measure is absolutely continuous for small α. In [SSU2], using some of the techniques developed in

this paper, we establish that the invariant measure is indeed absolutely continuous for a.e. α in some interval.

Here is a brief outline of the contents of the paper.

Section 2 contains preliminaries concerning infinite hyperbolic IFS, including properties of the pressure function.

Section 3 deals with families of infinite hyperbolic IFS depending on parameters. Its main result, which is of independent interest, is Theorem 3.1, which computes the Hausdorff dimension and Lebesgue measure of the limit set of a.e. infinite hyperbolic IFS with overlaps from a family satisfying a transversality condition.

The exceptional set of parameters, associated with Theorem 3.1, is analyzed in Section 4 where we estimate its local dimension from above.

In Section 5 we consider a single parabolic IFS. Following the approach of [MU2, MU3], we reduce the parabolic IFS to an infinite hyperbolic IFS. The limit sets of the parabolic and infinite hyperbolic systems differ in a countable set, so they have the same dimension. We prove that the unique zero of the pressure for the infinite hyperbolic IFS coincides with the least zero of the pressure for the parabolic IFS, even though the pressure functions for these systems differ (see Proposition 5.10(ii)).

In Section 6 we study families of parabolic IFS. Our main result, Theorem 6.1, computes the Hausdorff dimension and Lebesgue measure of the limit set of a.e. parabolic IFS with overlaps in a family satisfying a transversality condition. Moreover, applying the results of Section 4 we estimate the local dimension of the exceptional set, in the spirit of Proposition 1.1(ii).

Section 7 is devoted to examples. We consider two general classes of examples. The most difficult part in applying Theorem 6.1 is checking the transversality condition. In Propositions 7.1 and 7.2 we obtain effective sufficient conditions for transversality. We conclude with specific examples arising from continued fractions.

Notation. We write $B_\delta(\mathbf{t}_0)$ for the open ball of radius δ centered at $\mathbf{t}_0$ and $\mathcal{L}_d$ for the Lebesgue measure in $\mathbb{R}^d$. If μ is a measure we often write μA without parentheses. The symbol $\preceq$ means that the inequality holds up to an absolute multiplicative constant, and $\asymp$ means that both $\preceq$ and $\succeq$ are true.

2. Preliminaries

Let $X \subset \mathbb{R}$ be a closed interval. We consider a collection $\Psi = \{\psi_i\}_{i\in I}$ of continuous self-maps of X, where the set I may be finite or countable. We set $I^* := \bigcup_{n\geq 1} I^n$ and denote

by I^∞ the set of all infinite sequences of elements of I. If $\omega \in I^*$, then by $|\omega|$ we denote the length of ω. If $\omega \in I^* \cup I^\infty$ and $|\omega| > n$ then $\omega|_n = \omega_1\omega_2...\omega_n$ is the word consisting of the first n letters of ω and $\sigma^n\omega = \omega_{n+1}\omega_{n+2}...\omega_{|\omega|}$. The shift map $\sigma : I^\infty \to I^\infty$ sends an element $\{\omega_k\}_{k=1}^\infty$ to the element $\{\omega_{k+1}\}_{k=1}^\infty$. If $\omega \in I^n$, then by $\psi_\omega : X \to X$ we denote the composition $\psi_{\omega_1} \circ \psi_{\omega_2} \circ ... \circ \psi_{\omega_n}$. Notice that given $\omega \in I^\infty$, the sequence of compact sets $\{\psi_{\omega|_n}(X)\}_{n=1}^\infty$ is descending and therefore $\bigcap_{n\geq 1} \psi_{\omega|_n}(X) \neq \emptyset$. If for every $\omega \in I^\infty$ this intersection is a singleton, the collection Ψ is said to be a **topological IFS**. We can then define the map $\pi_\Psi : I^\infty \to X$ by setting

$$\{\pi_\Psi(\omega)\} = \bigcap_{n\geq 1} \psi_{\omega|_n}(X).$$

We call this map the **natural projection** induced by the topological IFS Ψ, and its range, the set $J_\Psi = \pi_\Psi(I^\infty)$, is called the **limit set** of Ψ. Thus, $\lim_{n\to\infty} \operatorname{diam}\Big(\psi_{\omega|_n}(X)\Big) = 0$ and therefore,

$$\pi_\Psi(\omega) = \lim_{n\to\infty} \psi_{\omega|_n}(x)$$

for every $x \in X$. We also have the following useful identity:

$$\pi_\Psi(\omega) = \psi_{\omega|_n}\left(\pi_\Psi(\sigma^n\omega)\right) \quad \text{for any } \ n \geq 1. \tag{2.1}$$

The limit set satisfies $J_\Psi = \bigcup_{i\in I} \psi_i(J_\Psi)$ but it need not be compact when I is infinite [MU1]. We call Ψ a **smooth IFS** if the following condition is satisfied.

SMOOTHNESS: there exists $\theta \in (0, 1]$ such that

$$\psi_i \in \mathcal{C}^{1+\theta}(X \to \operatorname{Int}(X)), \quad \text{and } \ \psi_i'(x) \neq 0 \quad \text{for all } \ x \in X \ \text{ and } \ i \in I. \tag{2.2}$$

Given $t \geq 0$ we define the pressure function $P_\Psi(t)$ by the formula

$$P_\Psi(t) = \lim_{n\to\infty} \frac{1}{n} \log Z_n(\Psi, t) \tag{2.3}$$

where $Z_n(\Psi, t) = \sum_{|\omega|=n} \| \psi_\omega' \|^t$ and $\| \cdot \|$ denotes the supremum norm on X. Observe that the limit in (2.3) really exists since the sequence $n \mapsto \log Z_n(\Psi, t)$ is subadditive. Comparing this with the usual definition of topological pressure (see e.g. [Bo1]) we see that $P_\Psi(t)$ equals the pressure of the shift map σ on I^∞ with the potential

$$\omega \mapsto t \log |\psi_{\omega_1}'(\pi(\sigma\omega))|.$$

We call a smooth IFS **hyperbolic** if the following conditions are satisfied:

HYPERBOLICITY: for some $\gamma \in (0, 1)$,

$$\| \psi_i' \| \leq \gamma < 1, \quad \text{for all } \ i \in I \tag{2.4}$$

(any map satisfying this property will be called hyperbolic), and

BOUNDED DISTORTION PROPERTY: there exists $K > 1$ such that for all $n \in \mathbb{N}$ and $\omega \in I^n$,

$$K^{-1} \le \frac{|\psi_\omega(x)|}{|\psi_\omega(y)|} \le K \quad \text{for all } \ x, y \in X. \tag{2.5}$$

Such Ψ is a one-dimensional example of **conformal IFS**, introduced (for infinite I) and explored in [MU1], where also the open set condition was assumed. It is well-known that (2.5) follows from (2.2) and (2.4) when I is finite, see e.g. [B1] (see [MU1] for more general sufficient conditions).

From now on throughout the Sections 2–4, unless otherwise stated, we assume that the smooth IFS Ψ is hyperbolic. Let

$$\Theta(\Psi) = \inf\{t \ge 0 : P_\Psi(t) < \infty\}. \tag{2.6}$$

The behavior of the pressure function is described in the following lemma.

Lemma 2.1. *The function $t \mapsto P_\Psi(t)$, for $t \in (\Theta(\Psi), \infty)$, is finite, strictly decreasing and continuous, and $\lim_{t\to\infty} P_\Psi(t) = -\infty$.*

Proof. The first statement is immediate from the definition of the number $\Theta = \Theta(\Psi)$. Now, given $t > \Theta$ and $s > 0$, we have by (2.4) for all $n \ge 1$:

$$Z_n(\Psi, t+s) = \sum_{|\omega|=n} \| \, (\psi_\omega)' \, \|^{t+s} \le \sum_{|\omega|=n} \| \, (\psi_\omega)' \, \|^{t} \, \gamma^{ns},$$

and therefore, $P_\Psi(t+s) \le s \log \gamma + P_\Psi(t) < P_\Psi(t)$. Thus, the function $t \mapsto P_\Psi(t)$ is strictly decreasing on $t \in (\Theta, \infty)$. Now, an application of Hölder's inequality shows that each function $t \mapsto Z_n(\Psi, t)$ is log convex. Therefore the function $t \mapsto P_\Psi(t)$, $t \in (\Theta, \infty)$, is convex and, consequently, continuous. By the definition of Θ we have $P_\Psi(\Theta + 1) < \infty$. Hence, for every $t > 0$ and every $n \ge 1$,

$$Z_n(\Psi, \Theta + 1 + t) = \sum_{|\omega|=n} \| \, \psi'_\omega \, \|^{\Theta+1+t} \le \sum_{|\omega|=n} \| \, \psi'_\omega \, \|^{1+\Theta} \cdot \| \, \psi'_\omega \, \|^{t} \le \gamma^{tn} \sum_{|\omega|=n} \| \, \psi'_\omega \, \|^{1+\Theta} \ .$$

Therefore, $P_\Psi(\Theta + 1 + t) \le t \log \gamma + P_\Psi(1 + \Theta)$ and hence $\lim_{s\to+\infty} P_\Psi(s) = -\infty$. The proof is complete. □

Definition 2.2. *Following* [MU1] *we call a hyperbolic system Ψ* **regular** *if*

$$\exists \, s(\Psi) \ge \Theta(\Psi) : \quad P_\Psi(s(\Psi)) = 0. \tag{2.7}$$

We denote by $\Xi_X(K, \gamma, \theta)$ the class of regular hyperbolic IFS on X.

In view of Lemma 2.1, if the number $s(\Psi)$ exists, then it is unique. Also, if $\#I < \infty$ then Ψ is regular since then $\Theta(\Psi) = 0$ and $P_\Psi(0) = \log(\#I) > 0$. The following lemma shows that $s(\Psi)$ is always an upper bound for the Hausdorff dimension of the limit set. The argument is well-known but we include it for completeness. We write $\mathcal{H}^\alpha(A) = \lim_{\delta\to 0} \mathcal{H}^\alpha_\delta(A)$ for the α-dimensional Hausdorff measure of a set A.

Lemma 2.3. *If $\Psi \in \Xi_X(K,\gamma,\theta)$, then $\dim_{\mathrm{H}}(J_\Psi) \le s(\Psi)$.*

Proof. Fix $\varepsilon, \delta > 0$ and take n_1 so large that $\gamma^{n_1} \le \delta$. By the definition of $s(\Psi)$ and Lemma 2.1 there exists $n_2 \ge n_1$ and $\eta > 0$ such that for all $n \ge n_2$,

$$\frac{1}{n}\log\left(\sum_{|\omega|=n} \| \psi'_\omega \|^{s(\Psi)+\varepsilon}\right) \le -\eta.$$

Hence, for all $n \ge n_2$,

$$\begin{aligned}\mathcal{H}^{s(\Psi)+\varepsilon}_\delta(J_\Psi) &\le \sum_{|\omega|=n} \operatorname{diam}\left(\psi_\omega(X)\right)^{s(\Psi)+\varepsilon} \\ &\le \operatorname{diam}(X)^{s(\Psi)+\varepsilon} \sum_{|\omega|=n} \| \psi'_\omega \|^{s(\Psi)+\varepsilon} \\ &\le \operatorname{diam}(X)^{s(\Psi)+\varepsilon} e^{-\eta n},\end{aligned}$$

and, consequently, $\mathcal{H}^{s(\Psi)+\varepsilon}_\delta(J_\Psi) = 0$. Thus, $\mathcal{H}^{s(\Psi)+\varepsilon}(J_\Psi) = 0$, and letting $\varepsilon \searrow 0$ we conclude that $\dim_{\mathrm{H}}(J_\Psi) \le s(\Psi)$. □

Given an IFS $\Psi = \{\psi_i : \ i \in I\}$ and $F \subset I$ let $\Psi_F = \{\psi_i : i \in F\}$. Denote $\mathit{Fin}(I) = \{F \subset I : \#F < \infty\}$. We are going to show that $s(\Psi)$ is the supremum of $s(\Psi_F)$ over $F \in \mathit{Fin}(I)$. Along the way we obtain an estimate on the speed of convergence in (2.3), which will be useful later.

Lemma 2.4. *Let $\Psi \in \Xi_X(K,\gamma,\theta)$. Then*

(i) *for every $t > 0$ and every $n \ge 1$,*

$$P_\Psi(t) \le \frac{1}{n}\log Z_n(\Psi,t) \le P_\Psi(t) + \frac{t\log K}{n}. \tag{2.8}$$

(ii) $s(\Psi) = \sup\{s(\Psi_F) : \ F \in \mathit{Fin}(I)\}$.

Proof. (i) The left-hand side inequality is immediate from the subadditivity of the sequence $n \mapsto \log Z_n(\Psi,t)$. In order to prove the right-hand side inequality, fix $n \ge 1$ and

consider an arbitrary integer $q \geq 1$. Then

$$\begin{aligned}\frac{1}{qn}\log Z_{qn}(\Psi,t) &= \frac{1}{qn}\log \sum_{\omega\in I^{qn}} \| \psi'_\omega \|^t \\ &\geq \frac{1}{qn}\log\left(K^{-qt}\sum_{\omega\in(I^n)^q}\prod_{i=1}^{q}\| \psi'_{\tau_i} \|^t\right) \\ &= \frac{-t\log K}{n} + \frac{1}{qn}\log\left(\sum_{\tau\in I^n}\| \psi'_\tau \|^t\right)^q \\ &= \frac{-t\log K}{n} + \frac{1}{n}\log Z_n(\Psi,t),\end{aligned}$$

where in the second displayed line we used (2.5) and $\omega = \tau_1\ldots\tau_q$, with $\tau_i \in I^n$. Letting $q\to\infty$ we obtain the right-hand side inequality of (2.8).

(ii) Since for every $t\geq 0$ and every $F\subset I$ we have $P_{\Psi_F}(t)\leq P_\Psi(t)$, the inequality $s(\Psi)\geq \sup\{s(\Psi_F):\ F\in \mathit{Fin}(I)\}$ is obvious. The opposite inequality will be deduced from (i). Fix an arbitrary $\Theta(\Psi) < t < s(\Psi)$. Then $0 < P_\Psi(t) < \infty$ and there exists $n\in\mathbb{N}$ so large that $P_\Psi(t) > \frac{2t\log K}{n}$. Fix such an n. Clearly, $Z_n(\Psi,t) = \sup\{Z_n(\Psi_F):\ F\in \mathit{Fin}(I)\}$, hence we can find $F\in \mathit{Fin}(I)$ satisfying

$$\frac{\log Z_n(\Psi_F,t)}{n} \geq \frac{\log Z_n(\Psi,t)}{n} - \frac{t\log K}{n} \geq P_\Psi(t) - \frac{t\log K}{n} > \frac{t\log K}{n}\,.$$

But now, applying (2.8) to $\Psi_F\in\Xi_X(K,\gamma,\theta)$ we obtain

$$P_{\Psi_F}(t) \geq \frac{\log Z_n(\Psi_F,t)}{n} - \frac{t\log K}{n} > 0$$

which implies that $t\leq s(\Psi_F)$. Thus, $s(\Psi)\leq \sup\{s(\Psi_G):\ G\in \mathit{Fin}(I)\}$, and the proof is complete. □

3. Families of hyperbolic IFS.

Let $X\subset\mathbb{R}$ be a compact interval and $U\subset\mathbb{R}^d$ an open set. Here we consider families of hyperbolic IFS $\Psi^{\mathbf{t}}\in\Xi_X(K,\gamma,\theta)$ depending on a parameter $\mathbf{t}\in U$. By $J_{\mathbf{t}}$ we denote the limit set of $\Psi^{\mathbf{t}}$ and by $\pi_{\mathbf{t}} = \pi_{\Psi^{\mathbf{t}}} : I^\infty \to J_{\mathbf{t}}$ we denote the natural projection introduced in Section 2. We need two conditions concerned with the dependence of the IFS on $\mathbf{t}$.

Distortion Continuity: for any $\eta>0$ there exists $\delta>0$ such that

$$\mathbf{t}_1,\mathbf{t}_2\in U,\ \|\mathbf{t}_1-\mathbf{t}_2\|\leq\delta \quad\Longrightarrow\quad \forall\omega\in I^*,\ \ e^{-|\omega|\eta}\leq\frac{\|(\psi^{\mathbf{t}_1}_\omega)'\|}{\|(\psi^{\mathbf{t}_2}_\omega)'\|}\leq e^{|\omega|\eta}\,. \tag{3.1}$$

TRANSVERSALITY CONDITION: for any ω and τ in I^∞ with $\omega_1 \neq \tau_1$, there exists a constant $C_1 = C_1(\omega_1, \tau_1)$ such that

$$\mathcal{L}_d\{\mathbf{t} \in U : \ |\pi_{\mathbf{t}}(\omega) - \pi_{\mathbf{t}}(\tau)| \leq r\} \leq C_1 r \quad \text{for all} \ \ r > 0. \tag{3.2}$$

We emphasize that C_1 depends only on ω_1 and τ_1. Thus, C_1 can be assumed independent of ω and τ if I is finite.

Now we can state the main result of this section. We write $s(\mathbf{t}) = s(\Psi^{\mathbf{t}})$.

Theorem 3.1. *Suppose that $\{\Psi^{\mathbf{t}}\}_{\mathbf{t}\in U}$ is a family of IFS in $\Xi_X(K, \gamma, \theta)$ satisfying (3.1) and (3.2). Then the function $\mathbf{t} \mapsto s(\mathbf{t})$ is continuous on U and*

(i) $\dim_{\mathrm{H}}(J_{\mathbf{t}}) = \min\{s(\mathbf{t}), 1\}$ *for Lebesgue-a.e.* $\mathbf{t} \in U$;

(ii) $\mathcal{L}_1(J_{\mathbf{t}}) > 0$ *for Lebesgue-a.e.* $\mathbf{t} \in U$ *such that* $s(\mathbf{t}) > 1$.

The rest of the section is devoted to the proof of this theorem. We begin with two lemmas which are easy consequences of (3.1) and (3.2).

Lemma 3.2. *Given $\varepsilon, a > 0$ define $\eta = \frac{-\varepsilon \log \gamma}{4a+\varepsilon}$ and take $\delta = \delta(\eta)$ coming from (3.1) ascribed to η. Then for all $\omega \in I^*$,*

$$\| \mathbf{t}_0 - \mathbf{t} \| < \delta \implies \| (\psi_\omega^{\mathbf{t}_0})' \|^{a+\frac{\varepsilon}{4}} \leq \| (\psi_\omega^{\mathbf{t}})' \|^a .$$

Proof. By (3.1) we have

$$\begin{aligned} \| (\psi_\omega^{\mathbf{t}_0})' \|^{a+\frac{\varepsilon}{4}} &\leq e^{|\omega|\eta(a+\frac{\varepsilon}{4})} \cdot \| (\psi_\omega^{\mathbf{t}})' \|^{a+\frac{\varepsilon}{4}} \\ &\leq e^{|\omega|\eta(a+\frac{\varepsilon}{4})} \cdot \gamma^{|\omega|\frac{\varepsilon}{4}} \cdot \| (\psi_\omega^{\mathbf{t}})' \|^a \\ &= e^{|\omega|(\eta(a+\frac{\varepsilon}{4})+\frac{\varepsilon}{4}\log\gamma)} \cdot \| (\psi_\omega^{\mathbf{t}})' \|^a = \| (\psi_\omega^{\mathbf{t}})' \|^a . \end{aligned}$$

The proof is complete. □

Lemma 3.3. *Suppose that the family $\Psi^{\mathbf{t}}$ satisfies (3.2). Then for every $0 < \alpha < 1$ and for all $\omega, \tau \in I^\infty$ with $\omega_1 \neq \tau_1$, there exists $C_2 = C_2(\alpha, \omega_1, \tau_1) > 0$ such that*

$$\int_U \frac{d\mathbf{t}}{|\pi_{\mathbf{t}}(\omega) - \pi_{\mathbf{t}}(\tau)|^\alpha} \leq C_2.$$

Proof. In view of (3.2), we can estimate as follows:

$$\begin{aligned}\int_U \frac{d\mathbf{t}}{|\pi_{\mathbf{t}}(\omega)-\pi_{\mathbf{t}}(\tau)|^\alpha} &= \int_0^\infty \mathcal{L}_d\left\{\mathbf{t}\in U:\ \frac{1}{|\pi_{\mathbf{t}}(\omega)-\pi_{\mathbf{t}}(\tau)|^\alpha}\ge x\right\}dx\\ &= \int_0^\infty \mathcal{L}_d\left\{\mathbf{t}\in U:\ |\pi_{\mathbf{t}}(\omega)-\pi_{\mathbf{t}}(\tau)|\le r\right\} r^{-\alpha-1}dr\\ &= \int_0^{|X|} \mathcal{L}_d\left\{\mathbf{t}\in U:\ |\pi_{\mathbf{t}}(\omega)-\pi_{\mathbf{t}}(\tau)|\le r\right\} r^{-\alpha-1}dr +\\ &+ \int_{|X|}^\infty \mathcal{L}_d\left\{\mathbf{t}\in U:\ |\pi_{\mathbf{t}}(\omega)-\pi_{\mathbf{t}}(\tau)|\le r\right\} r^{-\alpha-1}dr\\ &\le C_1(\omega_1,\tau_1)(1-\alpha)^{-1}|X|^{1-\alpha}+\mathcal{L}_d(U)\alpha^{-1}|X|^{-\alpha},\end{aligned}$$

and the lemma is proved. □

The next lemma is proved following the scheme of [SSo, Lemma 4.1(ii)]; it implies the continuity statement in Theorem 3.1.

Lemma 3.4. *If the family $\Psi^{\mathbf{t}}\in\Xi_X(K,\gamma,\theta)$, with $\mathbf{t}\in U$, satisfies (3.1), then the function $\mathbf{t}\mapsto s(\mathbf{t})$ is continuous on U.*

Proof. Fix an arbitrary $\Phi\in\Xi_X(K,\gamma,\theta)$. Then for every $t>\Theta(\Phi)$ and $u\ge 0$ we have

$$\begin{aligned}P_\Phi(t+u)-P_\Phi(t) &= \lim_{n\to\infty}\frac{1}{n}\left(\log\left(\sum_{|\omega|=n}\|\ \phi'_\omega\ \|^{t+u}\right)-\log\left(\sum_{|\omega|=n}\|\ \phi'_\omega\ \|^{t}\right)\right)\\ &\le \lim_{n\to\infty}\frac{1}{n}\left(\log\left(\gamma^{un}\sum_{|\omega|=n}\|\ \phi'_\omega\ \|^{t}\right)-\log\left(\sum_{|\omega|=n}\|\ \phi'_\omega\ \|^{t}\right)\right)\\ &= u\log\gamma.\end{aligned}$$

Therefore, for all $t>\Theta(\Phi)$ and $u>\Theta(\Phi)-t$,

$$|P_\Phi(t+u)-P_\Phi(t)|\ge |u|\cdot|\log\gamma|. \tag{3.3}$$

Recall that$P_{\Psi^{\mathbf{t}_i}}(s(\mathbf{t}_i))=0$ by the definition of $s(\mathbf{t}_i)=s(\Psi^{\mathbf{t}_i})$. Fix $\varepsilon>0$, consider $\delta>0$ produced by (3.1) with $\eta=\varepsilon$, and suppose that $\|\ \mathbf{t}_2-\mathbf{t}_1\ \|<\delta$. Then by (2.3) and (3.1),

$$|P_{\Psi^{\mathbf{t}_1}}(s(\mathbf{t}_2))| = |P_{\Psi^{\mathbf{t}_1}}(s(\mathbf{t}_2))-P_{\Psi^{\mathbf{t}_2}}(s(\mathbf{t}_2))| = \lim_{n\to\infty}\frac{1}{n}\log\left|\frac{\sum_{|\omega|=n}\|\ (\psi^{\mathbf{t}_1}_\omega)'\ \|^{s(\mathbf{t}_2)}}{\sum_{|\omega|=n}\|\ (\psi^{\mathbf{t}_2}_\omega)'\ \|^{s(\mathbf{t}_2)}}\right|\le s(\mathbf{t}_2)\varepsilon.$$

Therefore, $s(\mathbf{t}_2)>\Theta(\Psi^{\mathbf{t}_1})$ and in view of (3.3) we have

$$|s(\mathbf{t}_2)-s(\mathbf{t}_1)|\le\frac{1}{|\log\gamma|}\,|\mathrm{P}_{\Psi^{\mathbf{t}_1}}(s(\mathbf{t}_2))-\mathrm{P}_{\Psi^{\mathbf{t}_1}}(s(\mathbf{t}_1))| = \frac{1}{|\log\gamma|}\,|\mathrm{P}_{\Psi^{\mathbf{t}_1}}(s(\mathbf{t}_2))|\le\frac{s(\mathbf{t}_2)\varepsilon}{|\log\gamma|},$$

and the desired statement follows. □

Following [SSo], we now prove the main ingredient needed for the proof of Theorem 3.1.

Lemma 3.5. *Suppose that the family* $\{\Psi^{\mathbf{t}}\}_{\mathbf{t}\in U}$ *satisfies (3.1) and (3.2). Then*

(i) *for any* $\mathbf{t}_0 \in U$ *and any* $\varepsilon > 0$ *there exists* $\delta > 0$ *such that*

$$\dim_{\mathrm{H}}(J_{\mathbf{t}}) \geq \min\{s(\mathbf{t}_0), 1\} - \varepsilon \qquad \text{for } \mathcal{L}_d\text{-a.e. } \mathbf{t} \in B_\delta(\mathbf{t}_0).$$

(ii) *Suppose that* $s(\mathbf{t}_0) > 1 + \varepsilon$ *for some* $0 < \varepsilon < 1$. *Then there exists* $\delta > 0$ *such that*

$$\mathcal{L}_1(J_{\mathbf{t}}) > 0 \qquad \text{for } \mathcal{L}_d\text{-a.e. } \mathbf{t} \in B_\delta(\mathbf{t}_0).$$

Proof. Let $s = \min\{s(\mathbf{t}_0), 1\}$. By Lemma 2.4(ii), there exists a finite subset F of I such that $s(\Psi_F^{\mathbf{t}_0}) > s(\mathbf{t}_0) - \frac{\varepsilon}{2} \geq s - \frac{\varepsilon}{2}$. To simplify notation we set

$$\Psi = \{\psi_i\}_{i\in F} = \Psi_F^{\mathbf{t}_0}\,.$$

Consider the function $f:\ F^\infty \to \mathbb{R}$ defined by $f(\omega) = \log|\psi'_{\omega_1}(\pi_\Psi(\sigma\omega))|$. It follows from (2.2) that f is Hölder continuous, and (2.1) implies

$$\sum_{i=0}^{n-1} f(\sigma^i\omega) = \log|\psi'_{\omega|_n}(\pi_\Psi(\sigma^n\omega))| \quad \text{for all } \omega \in F^\infty.$$

Since $P_\Psi(s(\Psi)) = 0$, the theory of Gibbs states (see [Bo1], cf. [MU1] for a more general setting) produces a Borel probability shift-invariant measure μ on F^∞ such that for some constant $C_3 \geq 1$, all $\omega \in F^\infty$, and all $n \geq 1$,

$$\mu[\omega|_n] \in (C_3^{-1}, C_3)\,|\psi'_{\omega|_n}(\pi_\Psi(\sigma^n\omega))|^{s(\Psi)}\,.$$

Here $[\omega|_n]$ is the cylinder set of all sequences starting with $\omega_1 \ldots \omega_n$. The measure μ is called the Gibbs state for the potential $\omega \mapsto s(\Psi)f(\omega)$. Bounded distortion (2.5) implies that there exists a constant $C_4 \geq 1$ such that for all $\omega \in F^\infty$ and all $n \geq 1$,

$$\mu[\omega|_n] \in (C_4^{-1}, C_4)\ \|\ \psi'_{\omega|_n}\ \|^{s(\Psi)}\,. \tag{3.4}$$

Denote the product measure $\mu \times \mu$ by μ_2. First we prove part (i) of the lemma. By the potential-theoretic characterization of the Hausdorff dimension (see [F2, p.79]) it is enough to show that

$$R(\mathbf{t}) = \iint_{F^\infty\times F^\infty} \frac{d\mu_2(\omega,\tau)}{|\pi_{\mathbf{t}}(\omega) - \pi_{\mathbf{t}}(\tau)|^{s-\varepsilon}} < \infty \tag{3.5}$$

for a.e. $\mathbf{t} \in B_\delta(\mathbf{t}_0)$, where $\pi_{\mathbf{t}} = \pi_{\Psi^{\mathbf{t}}}$. Indeed, (3.5) means that the $(s-\varepsilon)$-energy of the "push-down" measure $\mu \circ \pi_{\mathbf{t}}^{-1}$, supported on the limit set $J_{\Psi_F^{\mathbf{t}}} \subset J_{\mathbf{t}}$, is finite.

Following the scheme of Kaufman [K] we prove that

$$\int_{B_\delta(\mathbf{t}_0)} R(\mathbf{t})d\mathbf{t} < \infty$$

where $\delta = \delta(\eta)$ comes from (3.1) and $\eta = \frac{-\varepsilon \log \gamma}{4(s(\Psi)-\frac{\varepsilon}{2})+\varepsilon}$. For $\rho \in F^n$ denote

$$A_\rho = \{(\omega, \tau) \in F^\infty \times F^\infty : \ \omega \wedge \tau = \rho\}$$

where $\omega \wedge \tau$ is the largest common initial segment of ω and τ. For $(\omega, \tau) \in A_\rho$ we have by (2.1), (2.5), and the Mean Value Theorem:

$$|\pi_{\mathbf{t}}(\omega) - \pi_{\mathbf{t}}(\tau)| = |(\psi_\rho^{\mathbf{t}})'(c)| \cdot |\pi_{\mathbf{t}}(\sigma^n\omega) - \pi_{\mathbf{t}}(\sigma^n\tau)| \geq K^{-1} \parallel (\psi_\rho^{\mathbf{t}})' \parallel \cdot |\pi_{\mathbf{t}}(\sigma^n\omega) - \pi_{\mathbf{t}}(\sigma^n\tau)| \tag{3.6}$$

By Lemma 3.2,

$$\parallel (\psi_\rho^{\mathbf{t}})' \parallel^{s-\epsilon} \geq \parallel (\psi_\rho^{\mathbf{t}})' \parallel^{s(\Psi)-\frac{\varepsilon}{2}} \geq \parallel \psi_\rho' \parallel^{s(\Psi)-\frac{\varepsilon}{4}} \quad \text{for } \ \mathbf{t} \in B_\delta(\mathbf{t}_0) \tag{3.7}$$

where $\psi_\rho = \psi_\rho^{\mathbf{t}_0}$. Now we can estimate using (3.6) and (3.7):

$$\begin{aligned}
\int_{B_\delta(\mathbf{t}_0)} R(\mathbf{t}) d\mathbf{t} &= \sum_{n\geq 0} \sum_{\rho \in F^n} \iint_{A_\rho} \left(\int_{B_\delta(\mathbf{t}_0)} \frac{d\mathbf{t}}{|\pi_{\mathbf{t}}(\omega) - \pi_{\mathbf{t}}(\tau)|^{s-\varepsilon}} \right) d\mu_2(\omega, \tau) \\
&\preceq \sum_{n\geq 0} \sum_{\rho \in F^n} \iint_{A_\rho} \left[\parallel \psi_\rho' \parallel^{-s(\Psi)+\frac{\varepsilon}{4}} \int_{B_\delta(\mathbf{t}_0)} \frac{d\mathbf{t}}{|\pi_{\mathbf{t}}(\sigma^n\omega) - \pi_{\mathbf{t}}(\sigma^n\tau)|^{s-\varepsilon}} \right] d\mu_2(\omega, \tau) \\
&\preceq \sum_{n\geq 0} \sum_{\rho \in F^n} \iint_{A_\rho} \parallel \psi_\rho' \parallel^{-s(\Psi)+\frac{\varepsilon}{4}} \, d\mu_2(\omega, \tau) \, .
\end{aligned}$$

In the last inequality we applied Lemma 3.3 which is possible because $(\sigma^n\omega)_1 \neq (\sigma^n\tau)_1$. The multiplicative constant depends on $(\sigma^n\omega)_1$ and $(\sigma^n\tau)_1$, but since the set F is finite this does not cause a problem. Next, applying (3.4) and (2.4) and denoting by $[\rho]$ the cylinder set of $\rho \in F^n$ we obtain:

$$\begin{aligned}
\int_{B_\delta(\mathbf{t}_0)} R(\mathbf{t}) d\mathbf{t} &\preceq \sum_{n\geq 0} \sum_{\rho \in F^n} \iint_{A_\rho} \frac{\parallel \psi_\rho' \parallel^{\frac{\varepsilon}{4}}}{\mu[\rho]} \, d\mu_2(\omega, \tau) \\
&\preceq \sum_{n\geq 0} \gamma^{\frac{n\varepsilon}{4}} \sum_{\rho \in F^n} \frac{\mu_2(A_\rho)}{\mu[\rho]} \\
&\leq \sum_{n\geq 0} \gamma^{\frac{n\varepsilon}{4}} \sum_{\rho \in F^n} \frac{(\mu[\rho])^2}{\mu[\rho]} = \sum_{n\geq 0} \gamma^{\frac{n\varepsilon}{4}} < \infty.
\end{aligned}$$

This concludes the proof of part (i).

(ii) Let $\eta = \frac{-\varepsilon \log \gamma}{4+\varepsilon}$ and determine $\delta = \delta(\eta)$ from (3.1). In view of Lemma 2.4, there exists a finite set $F \subset I$ such that $s(\Psi_F^{\mathbf{t}_0}) \geq 1 + \frac{\varepsilon}{2}$. We use the same set-up as in the proof of part (i) and let $\Psi = \Psi_F^{\mathbf{t}_0}$. Recall that μ is the Gibbs state for the potential $\omega \mapsto s(\Psi) \log |\psi'_{\omega_1}(\pi_\Psi(\sigma\omega))|$ satisfying (3.4). For every $\mathbf{t} \in B_\delta(\mathbf{t}_0)$ consider

$$\nu_{\mathbf{t}} = \mu \circ \pi_{\mathbf{t}}^{-1},$$

the push-down measure on the limit set $J_{\Psi^{\mathbf{t}}_F} \subset J_{\mathbf{t}}$. It is enough to show that $\nu_{\mathbf{t}}$ is absolutely continuous with respect to the Lebesgue measure $\mathcal{L}_1$ for a.e. $\mathbf{t} \in B_\delta(\mathbf{t}_0)$. We prove that

$$\mathcal{I} = \int_{B_\delta(\mathbf{t}_0)} \int_{\mathbb{R}} \underline{D}(\nu_{\mathbf{t}}, x)\, d\nu_{\mathbf{t}}\, d\mathbf{t} < \infty$$

where

$$\underline{D}(\nu_{\mathbf{t}}, x) = \liminf_{r \searrow 0} \frac{\nu_{\mathbf{t}}[x-r, x+r]}{2r}$$

is the lower density of the measure $\nu_{\mathbf{t}}$ at the point x. This will be sufficient since then for a.e. $\mathbf{t} \in B_\delta(\mathbf{t}_0)$ we will have $\underline{D}(\nu_{\mathbf{t}}, x) < \infty$ for $\nu_{\mathbf{t}}$-a.e. x and [M, Lemma 2.12] will imply that $\nu_{\mathbf{t}}$ is absolutely continuous. The argument below follows the scheme of [PSo1]. First we apply Fatou's Lemma to get

$$\mathcal{I} \le \liminf_{r \searrow 0} \int_{B_\delta(\mathbf{t}_0)} \int_{\mathbb{R}} \frac{\nu_{\mathbf{t}}[x-r, x+r]}{2r}\, d\nu_{\mathbf{t}}\, d\mathbf{t}. \tag{3.8}$$

Next we use the definition of $\nu_{\mathbf{t}}$ to change the variable, write $\nu_{\mathbf{t}}[x-r, x+r]$ as an integral of the indicator function, and change the variable once again to obtain

$$\int_{\mathbb{R}} \nu_{\mathbf{t}}[x-r, x+r]\, d\nu_{\mathbf{t}} = \iint_{F^\infty \times F^\infty} \mathbf{1}_{\{\omega \in F^\infty:\ |\pi_{\mathbf{t}}(\omega) - \pi_{\mathbf{t}}(\tau)| \le r\}}\, d\mu_2(\omega, \tau).$$

Substituting this into (3.8) and exchanging the order of integration yields

$$\begin{aligned}
\mathcal{I} &\le \liminf_{r \searrow 0} (2r)^{-1} \iint_{F^\infty \times F^\infty} \mathcal{L}_d\{\mathbf{t} \in B_\delta(\mathbf{t}_0) :\ |\pi_{\mathbf{t}}(\omega) - \pi_{\mathbf{t}}(\tau)| \le r\}\, d\mu_2(\omega, \tau) \\
&= \liminf_{r \searrow 0} (2r)^{-1} \sum_{n \ge 0} \sum_{\rho \in F^n} \iint_{A_\rho} \mathcal{L}_d\{\mathbf{t} \in B_\delta(\mathbf{t}_0) :\ |\pi_{\mathbf{t}}(\omega) - \pi_{\mathbf{t}}(\tau)| \le r\}\, d\mu_2(\omega, \tau)\,.
\end{aligned}$$

By (3.6), Lemma 3.2, and (3.2), we have for all $(\omega, \tau) \in A_\rho$:

$$\begin{aligned}
&\mathcal{L}_d\{\mathbf{t} \in B_\delta(\mathbf{t}_0) :\ |\pi_{\mathbf{t}}(\omega) - \pi_{\mathbf{t}}(\tau)| \le r\} \\
&\preceq \mathcal{L}_d \left\{ \mathbf{t} \in B_\delta(\mathbf{t}_0) :\ |\pi_{\mathbf{t}}(\sigma^n \omega) - \pi_{\mathbf{t}}(\sigma^n \tau)| \le \frac{Kr}{\| (\psi^{\mathbf{t}}_\rho)' \|} \right\} \\
&\le \mathcal{L}_d \left\{ \mathbf{t} \in B_\delta(\mathbf{t}_0) :\ |\pi_{\mathbf{t}}(\sigma^n \omega) - \pi_{\mathbf{t}}(\sigma^n \tau)| \le \frac{Kr}{\| \psi'_\rho \|^{1+\frac{\varepsilon}{4}}} \right\} \preceq r \parallel \psi'_\rho \parallel^{-1-\frac{\varepsilon}{4}}
\end{aligned}$$

Here we used again that the constant in (3.2) can be made independent of ω and τ, due to the fact that F is finite. Now we can estimate the integral $\mathcal{I}$ as follows:

$$\mathcal{I} \preceq \sum_{n \ge 0} \sum_{\rho \in F^n} \iint_{A_\rho} \parallel \psi'_\rho \parallel^{-1-\frac{\varepsilon}{4}}\, d\mu_2(\omega, \tau).$$

By (3.4) and (2.4),

$$\parallel \psi'_\rho \parallel^{-1-\frac{\varepsilon}{4}} \asymp (\mu[\rho])^{-\frac{1+\varepsilon/4}{s(\Psi)}} \le (\mu[\rho])^{-(1-\frac{\varepsilon}{12})} \preceq \gamma^{\frac{n\varepsilon}{12}} (\mu[\rho])^{-1},$$

since $s(\Psi) > 1 + \frac{\varepsilon}{2}$ and $(1 + \frac{\varepsilon}{4})/(1 + \frac{\varepsilon}{2}) < 1 - \frac{\varepsilon}{12}$ for $\varepsilon < 1$. Thus,

$$\mathcal{I} \preceq \sum_{n\geq 0} \gamma^{\frac{n\varepsilon}{12}} \sum_{\rho\in F^n} \frac{\mu_2(A_\rho)}{\mu[\rho]} \leq \sum_{n\geq 0} \gamma^{\frac{n\varepsilon}{12}} \sum_{\rho\in F^n} \mu[\rho] = \sum_{n\geq 0} \gamma^{\frac{n\varepsilon}{12}} < \infty.$$

The proof is complete. □

Proof of Theorem 3.1. By Lemma 2.3, the function $s(\mathbf{t}) = \min\{s(\Psi^{\mathbf{t}}), 1\}$ is an upper bound for the Hausdorff dimension of the limit set $J_{\mathbf{t}}$. So we just have to show that

$$\dim_{\mathrm{H}}(J_{\mathbf{t}}) \geq s(\mathbf{t})$$

for a.e. $\mathbf{t} \in U$. Suppose that this is not the case. Then we can find $\varepsilon > 0$ and $\mathbf{t}_0$, a density point of those $\mathbf{t}$ for which

$$\dim_{\mathrm{H}}(J_{\mathbf{t}}) < s(\mathbf{t}) - \varepsilon.$$

Then there exists $\delta_0 > 0$ such that for each $\delta < \delta_0$,

$$\mathcal{L}_d \left\{ \mathbf{t} \in B_\delta(\mathbf{t}_0) : \ \dim_{\mathrm{H}}(J_{\mathbf{t}}) < \min\{s(\mathbf{t}), 1\} - \varepsilon \right\} > 0. \tag{3.9}$$

However, by the continuity of the function $s(\mathbf{t})$ (see Lemma 3.4), if δ is small enough then $s(\mathbf{t}) < s(\mathbf{t}_0) + \frac{\varepsilon}{2}$ for all $\mathbf{t} \in B_\delta(\mathbf{t}_0)$. Thus, for all δ sufficiently small we obtain from (3.9) that

$$\mathcal{L}_d \left\{ \mathbf{t} \in B_\delta(\mathbf{t}_0) : \ \dim_{\mathrm{H}}(J_{\mathbf{t}}) < \min\{s(\mathbf{t}_0), 1\} - \frac{\varepsilon}{2} \right\} > 0.$$

This contradicts Lemma 3.5(i) and completes the proof of the first part of Theorem 3.1. The second part follows immediately from Lemma 3.5(ii). □

4. Exceptional parameters

In this section, following the scheme of Kaufman [K], we obtain an estimate from above for the local Hausdorff dimension of the set of exceptional parameters in Theorem 3.1(i). As in Section 3, we assume that $\{\Psi^{\mathbf{t}}\}_{\mathbf{t}\in U}$ is a family of IFS in $\Xi_X(K, \gamma, \theta)$ satisfying (3.1), but we will need the following stronger transversality condition which will be checked for all the examples that we consider. Denote by $N_r(F)$ the minimal number of balls of radius r needed to cover the set $F \subset \mathbb{R}^d$.

Strong Transversality Condition: for all ω and τ in I^∞ with $\omega_1 \neq \tau_1$, there exists a constant $C_1 = C_1(\omega_1, \tau_1)$ such that for all $r > 0$,

$$N_r\left(\{\mathbf{t} \in U : \ |\pi_{\mathbf{t}}(\omega) - \pi_{\mathbf{t}}(\tau)| \leq r\}\right) \leq C_1 r^{1-d}. \tag{4.1}$$

Of course, the strong transversality condition implies the transversality condition (3.2). In the same way as Lemma 3.3 we can prove the following.

Lemma 4.1. *Suppose that the family $\Psi^{\mathbf{t}}$ satisfies the strong transversality condition (4.1). Let m be a Borel probability measure in $\mathbb{R}^d$ such that $m(B_r(x)) \le Cr^u$ for some $C,\ u > 0$ and all $x \in \mathbb{R}^d,\ r > 0$. Then for every $\alpha < u - d + 1$ and for all $\omega, \tau \in \mathcal{A}^\infty$ with $\omega_1 \ne \tau_1$, there exists $C_2 = C_2(\alpha, \omega_1, \tau_1) > 0$ such that*

$$\int_U \frac{dm(\mathbf{t})}{|\pi_{\mathbf{t}}(\omega) - \pi_{\mathbf{t}}(\tau)|^\alpha} \le C_2.$$

In the sequel any measure with the properties required in Lemma 4.1 will be called a Frostman measure with exponent u. Next we prove the analog of Lemma 3.5(i).

Lemma 4.2. *Suppose that the family $\{\Psi^{\mathbf{t}}\}_{\mathbf{t}\in U}$ satisfies (3.1) and (4.1). Then for any $\mathbf{t}_0 \in U$ and any $\varepsilon > 0$ there exists $\delta = \delta(\mathbf{t}_0, \varepsilon) > 0$ such that if m is a Frostman measure on $B_\delta(\mathbf{t}_0)$ with exponent u, then*

$$\dim_{\mathrm{H}}(J_{\mathbf{t}}) \ge \min\{s(\mathbf{t}_0), u - d + 1\} - \varepsilon$$

for m-a.e. $\mathbf{t} \in B_\delta(\mathbf{t}_0)$.

Proof. We let $s = \min\{s(\mathbf{t}_0), u - d + 1\}$ and then repeat the proof of Lemma 3.5(i) almost word by word. The only change is that now we prove that $\int_{B_\delta(\mathbf{t}_0)} R(\mathbf{t}) dm(\mathbf{t}) < \infty$ using Lemma 4.1 in the place where Lemma 3.3 was used. □

Now we can prove the main result of this section.

Theorem 4.3. *Suppose that the d-parameter family of IFS $\{\Psi^{\mathbf{t}}\}_{\mathbf{t}\in U}$ satisfies (3.1) and (4.1). If G is an arbitrary subset of U, then for every $\xi > 0$ we have*

$$\dim_{\mathrm{H}}\left(\left\{\mathbf{t} \in G :\ \dim_{\mathrm{H}}(J_t) < \min\{\xi, s(\mathbf{t})\}\right\}\right) \le \min\{\xi, \sup_G s(\mathbf{t})\} + d - 1.$$

Proof. Denote $\kappa := \min\{\xi, \sup_G s(\mathbf{t})\} + d - 1$. By the countable stability of the Hausdorff dimension, it is enough to prove that for all $n \in \mathbb{N}$,

$$\dim_{\mathrm{H}}\left(\left\{\mathbf{t} \in G :\ \dim_{\mathrm{H}}(J_t) < \min\{\xi, s(\mathbf{t})\} - \frac{1}{n}\right\}\right) \le \kappa.$$

Fix n and observe that it suffices to show that for all $\mathbf{t}_0$ in G there exists $\delta = \delta(\mathbf{t}_0)$ such that

$$\dim_{\mathrm{H}}\left(\left\{\mathbf{t} \in B_\delta(\mathbf{t}_0) :\ \dim_{\mathrm{H}}(J_t) < \min\{\xi, s(\mathbf{t})\} - \frac{1}{n}\right\}\right) \le \kappa$$

(just use that any cover of G contains a countable subcover and again the countable stability of the Hausdorff dimension). To establish our claim, suppose that it is false. Then there exists $\mathbf{t}_0$ such that for all $\delta > 0$

$$\dim_{\mathrm{H}}\left(\left\{\mathbf{t} \in B_\delta(\mathbf{t}_0) :\ \dim_{\mathrm{H}}(J_t) < \min\{\xi, s(\mathbf{t})\} - \frac{1}{n}\right\}\right) > \kappa.$$

Choose $\delta > 0$ so small that the statement of Lemma 4.2 holds with $\varepsilon = \frac{1}{2n}$ and $|s(\mathbf{t}) - s(\mathbf{t}_0)| < \frac{1}{2n}$ for all $\mathbf{t} \in B_\delta(\mathbf{t}_0)$ (by the continuity of $s(\mathbf{t})$). Then

$$\Big\{\mathbf{t} \in B_\delta(\mathbf{t}_0) :\ \dim_{\mathrm{H}}(J_t) < \min\{\xi, s(\mathbf{t})\} - \frac{1}{n}\Big\}$$
$$\subset \Big\{\mathbf{t} \in B_\delta(\mathbf{t}_0) :\ \dim_{\mathrm{H}}(J_t) < \min\{\xi, s(\mathbf{t}_0)\} - \frac{1}{2n}\Big\} =: E,$$

hence $\dim_{\mathrm{H}}(E) > \kappa$. By Frostman's Lemma (see [M, Th.8.8]), there is a Frostman measure m on the set E with exponent $u = \kappa$. By Lemma 4.2, for m-a.e. $\mathbf{t}$ we have

$$\dim_{\mathrm{H}}(J_\mathbf{t}) \geq \min\{s(\mathbf{t}_0), \kappa - d + 1\} - \frac{1}{2n} = \min\Big\{s(\mathbf{t}_0), \min\{\xi, \sup_G s(\mathbf{t})\}\Big\} - \frac{1}{2n}.$$

This is a contradiction since for all $\mathbf{t} \in E$ we have $\dim_{\mathrm{H}}(J_t) < \min\{\xi, s(\mathbf{t}_0)\} - \frac{1}{2n}$ and

$$\min\{\xi, s(\mathbf{t}_0)\} \leq \min\Big\{s(\mathbf{t}_0), \min\{\xi, \sup_G s(\mathbf{t})\}\Big\}.$$

□

Since the function $\mathbf{t} \mapsto s(\mathbf{t})$, $\mathbf{t} \in U$, is continuous, as an immediate consequence of Theorem 4.3 we get the following estimate for the local dimension of the exceptional set.

Corollary 4.4. *For every $\mathbf{t}_0 \in U$ we have*

$$\lim_{r\to 0} \dim_{\mathrm{H}}\Big(\{t \in B_r(\mathbf{t}_0) :\ \dim_{\mathrm{H}}(J_\mathbf{t}) < \min\{\xi, s(\mathbf{t})\}\}\Big) \leq \min\{\xi, s(\mathbf{t}_0)\} + d - 1.$$

5. Parabolic IFS.

Let $X \subset \mathbb{R}$ be a compact interval. We say that a $\mathcal{C}^{1+\theta}$ map $\phi : X \to X$ is **parabolic** if the following requirements are fulfilled:

- there is only one point $v \in X$ such that $\phi(v) = v$;
- $|\phi'(v)| = 1$ and $0 < |\phi'(x)| < 1$ for all $x \in X \setminus \{v\}$.
- There exists $L_1 \geq 1$ and $\beta = \beta(\phi) < \theta/(1-\theta)$ ($= \infty$ if $\theta = 1$) such that

$$L_1^{-1} \leq \liminf_{x\to v} \frac{|\,|\phi'(x)| - 1|}{|x - v|^\beta} \leq \limsup_{x\to v} \frac{|\,|\phi'(x)| - 1|}{|x - v|^\beta} \leq L_1.$$

At the beginning of this section we state some useful properties of a single parabolic map. They are very similar to [U, Lemmas 2.1–2.3]. First, integrating the partial sums of the series $\sum_{n=1}^\infty |(\phi^n)'(x)|$ we get the following.

Lemma 5.1. *For all $x \in \phi(X) \setminus \{v\}$ we have*

$$\frac{|x - v|}{|\phi^{-1}(x) - x|} \leq \sum_{n=1}^\infty |(\phi^n)'(x)| \leq \frac{|\phi(x) - v|}{|x - \phi(x)|}$$

Sending a sufficiently small neighborhood of v to infinity via the mapping $x \mapsto 1/(x-v)$, one can easily prove the following two local results.

Lemma 5.2. *For every neighborhood V of v there exists $L_2(V) \geq 1$ such that for all $x \in X \setminus V$ and all $n \geq 1$,*

$$\frac{1}{L_2(V)} \leq |\phi^n(x) - v| \cdot n^{1/\beta} \leq L_2(V).$$

Lemma 5.3. *For every neighborhood V of v there exists $L_3(V) \geq 1$ such that for all $x \in X \setminus V$ and all $n \geq 1$*

$$\frac{1}{L_3(V)} \leq |(\phi^n)'(x)| \cdot n^{\frac{\beta+1}{\beta}} \leq L_3(V).$$

Since $\beta < \theta/(1-\theta)$, the following is immediate from Lemma 5.3.

Corollary 5.4. *For every neighborhood V of v there exists $L_4(V) < \infty$ such that*

$$\sum_{n=1}^{\infty} \| (\phi^n)' \|^{\theta}_{X\setminus V} < L_4(V).$$

where $\| \cdot \|_{X\setminus V}$ denotes the sup-norm on $X \setminus V$.

Turning now our attention to iterated function systems we recall that a $\mathcal{C}^{1+\theta}$ map ϕ is hyperbolic if $0 < |\phi'(x)| < 1$ for all $x \in X$.

Definition 5.5. *Let $\Phi = \{\phi_1, \ldots, \phi_k\}$ be a collection of $\mathcal{C}^{1+\theta}$ functions on a closed interval $X \subset \mathbb{R}$ such that ϕ_k is parabolic with the fixed point v and the other functions are hyperbolic. We write $\Phi \in \Gamma_X(\theta)$ if, in addition,*

$$\phi_i(X) \subset \mathrm{Int}(X) \setminus \{v\} \quad \text{for all } \ i \leq k-1. \tag{5.1}$$

Remark. We consider IFS with just one parabolic element. The case of more than one parabolic function can also be handled, but at the cost of additional technical complications.

Let $\mathcal{A} = \{1, \ldots, k\}$, $\mathcal{A}^* := \bigcup_{n\geq 1} \mathcal{A}^n$, and suppose that

$$\max\{\| \phi_i' \| : \ i \leq k-1\} \leq \gamma < 1. \tag{5.2}$$

Lemma 5.6. *An IFS $\Phi \in \Gamma_X(\theta)$ is a topological IFS.*

Proof. All we need to show is that the intersections $\bigcap_{n\geq 1}\phi_{\omega|_n}(X)$ are singletons for all $\omega\in\mathcal{A}^\infty$. By (5.2),

$$\lim_{n\to\infty}\operatorname{diam}(\phi_{\omega|_n}(X))\to 0 \tag{5.3}$$

if the sequence ω has infinitely many symbols other than k. The remaining possibility is $\omega=wk^\infty$ for some $w\in\mathcal{A}^*$ but then (5.3) is still true since $\bigcap_{n\geq 1}\phi_{k^n}(X)=\{v\}$. □

Now, following [MU2, MU3], given a parabolic IFS $\Phi\in\Gamma_X(\theta)$, consider an associated infinite IFS

$$\Phi_*=\{\phi_k^n\phi_i:\ n\geq 0,\ i\leq k-1\}.$$

We also write $\Phi_*=\{\phi_b^*\}_{b\in I}$ where $\phi_b^*=\phi_k^n\phi_i$ and

$$I=\{b=(n,i):\ n\geq 0,\ i\leq k-1\}.$$

The following properties of Φ_* are immediate from the definitions.

Corollary 5.7. *Let $\Phi\in\Gamma_X(\theta)$. Then $\Phi_*=\{\phi_b^*\}_{b\in I}$ satisfies*

(i) *$\phi_b^*\in\mathcal{C}^{1+\theta}(X\to\operatorname{Int}(X))$ for all $b\in I$;*

(ii) *$0<\|\ (\phi_b^*)'\ \|\leq\gamma<1$ for all $b\in I$;*

(iii) *$J_\Phi=J_{\Phi_*}\cup\{\phi_w(v)\}_{w\in\mathcal{A}^*}$ so $\dim_{\mathrm{H}}(J_\Phi)=\dim_{\mathrm{H}}(J_{\Phi_*})$.*

Thus there is an infinite hyperbolic IFS Φ_* associated with a finite parabolic IFS Φ. This idea essentially goes back to Schweiger's "jump transformation" [Sc]. Our next goal is to show that $\Phi\in\Gamma_X(\theta)$ implies $\Phi_*\in\Xi_X(K,\gamma,\theta)$ for some K, see Definition 2.2. To achieve this, two more properties have to be verified: the bounded distortion property (2.5) and regularity (2.7). Bounded distortion properties for parabolic IFS (without overlaps) were investigated in [U]. Here a different version is needed but the approach is similar.

In the next section we study families of parabolic IFS, and it will be very important to know exactly what the various constants depend on. Therefore, we introduce the following notation. Let $\Phi\in\Gamma_X(\theta)$. We write

$$\Phi\in\Gamma_X(\theta,V,\gamma,u,M) \tag{5.4}$$

if V is a connected open neighborhood of the parabolic point v such that

$$V\cap\bigcup_{i=1}^{k-1}\phi_i(X)=\emptyset, \tag{5.5}$$

$$\max\{\|\ (\phi_i)'\ \|,\ i\leq k-1\}\leq\gamma\in(0,1), \tag{5.6}$$

$$\min\{|\phi_i'(x)|,\ x\in X,\ i\leq k\}\geq u\in(0,1), \tag{5.7}$$

and

$$\| \Phi' \|_\theta := \max_{i\le k} \sup\{|\phi_i'(x) - \phi_i'(y)| \cdot |x-y|^{-\theta}\} \le M. \tag{5.8}$$

By Definition 5.5, every $\Phi \in \Gamma_X(\theta)$ belongs to some $\Gamma_X(\theta, V, \gamma, u, M)$.

The next lemma will also be useful when we consider families of parabolic IFS. Recall that $\| \cdot \|_{X\setminus V}$ denotes the supremum norm on $X \setminus V$.

Lemma 5.8. *There exist constants $L_5 = L_5(X, \theta, V, \gamma, u, M) > 1$ and $L_6 = L_6(\theta, V, \gamma) > 0$ such that for every $\Phi \in \Gamma_X(\theta, V, \gamma, u, M)$, all $\omega \in \mathcal{A}^*$, and all $x, y \in X \setminus V$,*

$$L_5^{-|x-y|^\theta} \le \frac{|\phi_\omega'(y)|}{|\phi_\omega'(x)|} \le L_5^{|x-y|^\theta} \tag{5.9}$$

and

$$\sum_{j=0}^{|\omega|-1} \| \phi'_{\sigma^j\omega} \|_{X\setminus V}^\theta \le L_6. \tag{5.10}$$

Proof. We start with (5.10). Every $\tau \in \mathcal{A}^*$ can be written as

$$\tau = k^{r_1} j_1 k^{r_2} j_2 \dots k^{r_l} j_l k^{r_{l+1}} \tag{5.11}$$

where $l \ge 0$ and $r_p \ge 0, j_p \le k-1$ for $p \le l$. When $l = 0$ the equation (5.11) becomes $\tau = k^{r_1}$. One readily estimates using (5.5) and (5.6):

$$\| \phi'_\tau \|_{X\setminus V} \le \gamma^l \prod_{p=1}^{l+1} \| (\phi_k^{r_p})' \|_{X\setminus V} \ .$$

Applying this inequality to $\tau = \sigma^j\omega$ and summing over j we obtain

$$\begin{aligned}\sum_{j=0}^{|\omega|-1} \| \phi'_{\sigma^j\omega} \|_{X\setminus V}^\theta \le{}& \gamma^{l\theta} \cdot \left(\sum_{j=1}^{r_1} \| (\phi_k^j)' \|_{X\setminus V}^\theta + 1 \right) \cdot \prod_{p=2}^{l+1} \| (\phi_k^{r_p})' \|_{X\setminus V}^\theta \\ &+ \gamma^{(l-1)\theta} \cdot \left(\sum_{j=1}^{r_2} \| (\phi_k^j)' \|_{X\setminus V}^\theta + 1 \right) \cdot \prod_{p=3}^{l+1} \| (\phi_k^{r_p})' \|_{X\setminus V}^\theta \\ &+ \dots + \gamma^\theta \cdot \left(\sum_{j=1}^{r_l} \| (\phi_k^j)' \|_{X\setminus V}^\theta + 1 \right) \cdot \| (\phi_k^{r_{l+1}})' \| + \sum_{j=1}^{r_{l+1}} \| (\phi_k^j)' \|_{X\setminus V}^\theta \end{aligned}$$

Applying Corollary 5.4 for $\phi = \phi_k$ and using that $\| (\phi_k^r)' \|_{X\setminus V} \le 1$ for all r, we obtain

$$\sum_{j=0}^{|\omega|-1} \| \phi'_{\sigma^j\omega} \|_{X\setminus V}^\theta \le (L_4(V) + 1) \sum_{i=0}^{\infty} \gamma^{i\theta} =: L_6 = L_6(\theta, V, \gamma) < \infty,$$

and (5.10) is proved.

Now we turn to (5.9); clearly, it suffices to prove the right-hand side inequality. Suppose first that x and y belong to the same connected component of $X \setminus V$. Then using the Mean Value Theorem we conclude that for every $0 \le j \le |\omega|$ there exists $c_j \in X \setminus V$ such that $|\phi_{\sigma^j\omega}(y) - \phi_{\sigma^j\omega}(x)| = |\phi'_{\sigma^j\omega}(c_j)| \cdot |y - x|$. We have for all $1 \le i \le |\omega|$ by (5.7):

$$\begin{aligned}\left|\log|\phi'_{\omega_j}(\phi_{\sigma^j\omega}(y))| - \log|\phi'_{\omega_j}(\phi_{\sigma^j\omega}(x)|\right| &\le \left|\| \phi'_{\omega_j}(\phi_{\sigma^j\omega}(y)|\frac{1}{u} - |\phi'_{\omega_j}(\phi_{\sigma^j\omega}(x)|\right| \\ &\le \frac{\| \Phi' \|_\theta}{u} |\phi_{\sigma^j\omega}(y) - \phi_{\sigma^j\omega}(x)|^\theta \\ &= \frac{\| \Phi' \|_\theta}{u} |\phi'_{\sigma^j\omega}(c_j)|^\theta |y - x|^\theta.\end{aligned}$$

Since $c_j \in X \setminus V$, we obtain

$$\begin{aligned}|\log|\phi'_\omega(y)| - \log|\phi'_\omega(x)|| &= \left|\sum_{j=1}^{|\omega|} \log|\phi'_{\omega_j}(\phi_{\sigma^j\omega}(y))| - \sum_{j=1}^{|\omega|} \log|\phi'_{\omega_j}(\phi_{\sigma^j\omega}(x))|\right| \\ &\le \sum_{j=1}^{|\omega|} \left|\log|\phi'_{\omega_j}(\phi_{\sigma^j\omega}(y))| - \log|\phi'_{\omega_j}(\phi_{\sigma^j\omega}(x))|\right| \\ &\le \sum_{j=1}^{|\omega|} \frac{\| \Phi' \|_\theta}{u} |\phi'_{\sigma^j\omega}(c_j)|^\theta |y - x|^\theta \\ &\le \frac{\| \Phi' \|_\theta}{u} L_6 |y - x|^\theta \le \frac{M}{u} L_6 |y - x|^\theta.\end{aligned} \tag{5.12}$$

In the last displayed line we used (5.10) and (5.8). This completes the proof of (5.9) when x and y are in the same connected component of $X \setminus V$. Now suppose that they are in different components. Since then $|y - x| \ge \operatorname{diam}(V)$, it suffices to show the existence of a constant $L_7 = L_7(X, \theta, V, \gamma, u, M) \ge 1$ such that $|\phi'_\omega(y)| \le L_7 \cdot |\phi'_\omega(x)|$ for all $\omega \in \mathcal{A}^*$. To this end, suppose $\omega = \tau k^n$ where $n \ge 0$ and $\tau_{|\tau|} \ne k$. Let $|\tau| = l$. Observe that the points $\phi_{\tau_l k^n}(x)$ and $\phi_{\tau_l k^n}(y)$ belong to $\phi_{\tau_l}(X)$ and hence are in the same connected component of $X \setminus V$. Thus, in view of Lemma 5.3 and (5.12) we get

$$\frac{|\phi'_\omega(y)|}{|\phi'_\omega(x)|} = \frac{|\phi'_{\tau|_{l-1}}(\phi_{\tau_l k^n}(y))| \cdot |\phi'_{\tau_l}(\phi_{k^n}(y))| \cdot |\phi'_{k^n}(y)|}{|\phi'_{\tau|_{l-1}}(\phi_{\tau_l k^n}(x))| \cdot |\phi'_{\tau_l}(\phi_{k^n}(x))| \cdot |\phi'_{k^n}(x)|}$$

$$\le \exp\left(\frac{M}{u} L_6 |y - x|^\theta\right) \frac{\gamma}{u} L_3(V)^2,$$

and it remains to note that $|y - x| \le \operatorname{diam}(X)$. The only possibility left is $\omega = k^n$ but then $|\phi'_{k^n}(y)| \le L_3(V)^2 |\phi'_{k^n}(x)|$ by Lemma 5.3. The proof is complete. □

Corollary 5.9. *There exists a constant $K_1 = K_1(X, \theta, V, \gamma, u, M) > 1$ such that for every $\Phi \in \Gamma_X(\theta, V, \gamma, u, M)$, the associated infinite IFS Φ_* satisfies (2.5) with $K = K_1$. More precisely, for all $\tau \in I^*$ and all $x, y \in X$,*

$$K_1^{-|y-x|^\theta} \leq \frac{|(\phi_\tau^*)'(y)|}{|(\phi_\tau^*)'(x)|} \leq K_1^{|y-x|^\theta}$$

Proof. It is enough to prove only the right-hand side of this formula. By the definition of the system Φ_* we can write $\phi_\tau^* = \phi_\omega \phi_i$, where $\omega \in \mathcal{A}^*$ and $i \in \{1, 2, \ldots, k-1\}$. Then, applying (5.9) and (5.7) we can estimate as follows:

$$\begin{aligned}
\frac{|(\phi_\tau^*)'(y)|}{|(\phi_\tau^*)'(x)|} = \frac{|(\phi_\omega \circ \phi_i)'(y)|}{|(\phi_\omega \circ \phi_i)'(x)|} &= \frac{|\phi_\omega'(\phi_i(y))|}{|\phi_\omega'(\phi_i(x))|} \cdot \frac{|\phi_i'(y)|}{|\phi_i'(x)|} \\
&\leq L_5^{|\phi_i(y)-\phi_i(x)|^\theta} \exp\left(\frac{|\phi_i'(y)|}{|\phi_i'(x)|} - 1\right) \\
&\leq L_5^{|y-x|^\theta} \exp\left(\frac{\| \Phi' \|_\theta}{u} |y-x|^\theta\right) \\
&\leq L_5^{|y-x|^\theta} \exp\left(\frac{M}{u} |y-x|^\theta\right),
\end{aligned}$$

and the proof is finished. □

The pressure function $P_\Phi(t)$ for the parabolic IFS Φ is defined by

$$P_\Phi(t) = \lim_{n\to\infty} \frac{1}{n} \sum_{|\omega|=n} \| \phi_\omega' \|^t,$$

in accordance with (2.3). Observe that $P_\Phi(0) = \log k$. In contrast with the hyperbolic case, $P_\Phi(t) \geq 0$ for all $t > 0$ since $\| \phi_{k^n}' \| = 1$ for all n. It is easy to see that $P_\Phi(t)$ is non-increasing and continuous. Thus, $P_\Phi(s) = 0$ implies $P_\Phi(t) = 0$ for all $t > s$. Denote

$$s(\Phi) = \min\{t > 0 : \ P_\Phi(t) = 0\}.$$

If the pressure function has no zeros, we let $s(\Phi) = \infty$, however, the next proposition implies that $s(\Phi) < \infty$ for any $\Phi \in \Gamma_X(\theta)$.

Proposition 5.10. *Let $\Phi \in \Gamma_X(\theta)$ and let Φ_* be the associated infinite hyperbolic IFS. Then*

(i) *$\Theta(\Phi_*) = \frac{\beta}{\beta+1}$ and Φ_* is regular, i.e. there exists a unique $s(\Phi_*) > \Theta(\Phi_*)$ such that $P_{\Phi_*}(s(\Phi_*)) = 0$;*

(ii) *$s(\Phi_*) = s(\Phi)$.*

Before the proof, we point out the following.

Corollary 5.11. *If $\Phi \in \Gamma_X(\theta)$ then*

(i) *$\Phi_* \in \Xi_X(K_1, \gamma, \theta)$ where K_1 is from Corollary 5.9;*

(ii) $\dim_{\mathrm{H}}(J_\Phi) \le s(\Phi)$.

Proof: (i) follows from Proposition 5.10(i) and Definition 2.2(iii).

(ii) follows from Proposition 5.10(ii), Lemma 2.3, and Corollary 5.7. □

Proof of Proposition 5.10. (i) First we compute $\Theta(\Phi_*)$, see (2.6) for the definition. It follows from (5.1) and Lemma 5.3 that there exists a constant $K_2 > 1$ such that for all $x \in X$ and $b = (l, i) \in I$,

$$|(\phi_b^*)'(x)| = |(\phi_k^l \phi_i)'(x)| \in (K_2^{-1}, K_2)\,(l+1)^{-\frac{\beta+1}{\beta}}. \tag{5.13}$$

Therefore,

$$|(\phi_b^*)'(x)| \ge K_2^{-2} \parallel (\phi_b^*)' \parallel \qquad \text{for all } \ x \in X. \tag{5.14}$$

Let $t > \frac{\beta}{\beta+1}$. We have by (5.14), writing $\phi_\tau^* = \phi_{\tau_1}^* \dots \phi_{\tau_n}^*$:

$$\begin{aligned} Z_n(\Phi_*, t) = \sum_{|\tau|=n} \parallel (\phi_\tau^*)' \parallel^t &\in (K_2^{-2tn}, 1) \sum_{|\tau|=n} \parallel (\phi_{\tau_1}^*)' \parallel^t \dots \parallel (\phi_{\tau_n}^*)' \parallel^t \\ &= (K_2^{-2tn}, 1) \left(\sum_{b \in I} \parallel (\phi_b^*)' \parallel^t \right)^n \\ &= (K_2^{-2tn}, 1) \left(\sum_{i \le k-1} \sum_{l \ge 0} \parallel (\phi_k^l \phi_i)' \parallel^t \right)^n. \end{aligned}$$

Next we use (5.13) to get

$$\begin{aligned} Z_n(\Phi_*, t) &\in (K_2^{-3tn}, K_2^{tn}) \left(\sum_{i \le k-1} \sum_{l \ge 0} (l+1)^{-\frac{\beta+1}{\beta} t} \right)^n \\ &= \left(K_2^{-3tn} (k-1)^n, K_2^{tn} (k-1)^n \right) \left(\sum_{l \ge 0} (l+1)^{-\frac{\beta+1}{\beta} t} \right)^n. \end{aligned}$$

Since $P_{\Phi_*}(t) = \lim_{n\to\infty} \frac{1}{n} \log Z_n(\Phi_*, t)$ we obtain

$$\log(k-1) - 3t \log K_2 \le P_{\Phi_*}(t) - \sum_{l \ge 0} (l+1)^{-\frac{\beta+1}{\beta} t} \le \log(k-1) + t \log K_2.$$

It follows that $P_{\Phi_*}(t) < \infty$ for $t > \frac{\beta}{\beta+1}$ and $\lim_{t \searrow \frac{\beta}{\beta+1}} P_{\Phi_*}(t) = +\infty$. Thus, $\Theta(\Phi_*) = \frac{\beta}{\beta+1}$. Since $P_{\Phi_*}(t)$ is positive and finite for some t, we conclude from Lemma 2.1 that there exists a unique solution $s(\Phi_*)$ for the Bowen's equation $P_{\Phi_*}(t) = 0$. Part (i) is proved.

(ii) First we observe that

$$s(\Phi_*) = s' := \inf\left\{ t > 0 : \sum_{n\geq 1} Z_n(\Phi_*, t) < \infty \right\}. \tag{5.15}$$

Indeed, since $P_{\Phi_*}(t)$ is strictly decreasing, $s(\Phi_*) = \inf\{t > 0 : P_{\Phi_*}(t) < 0\}$. If $P_{\Phi_*}(t) < -\varepsilon < 0$ then $Z_n(\Phi_*, t) \leq Ce^{-\varepsilon n}$ as $n \to \infty$ and so $t \geq s'$. Thus, $s(\Phi_*) \geq s'$. On the other hand, if $\sum_{n\geq 1} Z_n(\Phi_*, t) < \infty$ then $Z_n(\Phi_*, t) \to 0$ hence $P_{\Phi_*}(t) \leq 0$. Therefore, $s(\Phi_*) \leq s'$ and (5.15) is proved.

Next we demonstrate that for all $t > 0$,

$$Z_n(\Phi, t) \leq 1 + \sum_{m\leq n} Z_m(\Phi_*, t). \tag{5.16}$$

Indeed, every $\omega \in \mathcal{A}^n$ can be written as $\omega = k^n$ or

$$\omega = k^{r_1} i_1 k^{r_2} i_2 \ldots k^{r_l} i_l k^{r_{l+1}}$$

where $l \geq 1$ and $i_p \leq k-1,\ r_p \geq 0$, for $p \leq l+1$. Thus, either $\phi_\omega = \phi_k^n$ or $\phi_\omega = \phi_\tau^* \phi_k^{r_{l+1}}$ for some $\tau \in I^*$. In the latter case, $\| \phi'_\omega \| \leq \| (\phi_\tau^*)' \|$ and $|\tau| = l \leq n$. Moreover, every map ϕ_τ^* with $|\tau| \leq n$ occurs in this procedure at most once. The estimate (5.16) follows by noting that $\| \phi'_{k^n} \| = 1$.

Now we can show that $s(\Phi_*) \geq s(\Phi) = \min\{t > 0 : P_\Phi(t) = 0\}$. In fact, if $\sum_{n\geq 1} Z_n(\Phi_*, r) < \infty$ then $Z_n(\Phi, r)$ is bounded by (5.16) and hence

$$P_\Phi(r) = \lim_{n\to\infty} \frac{1}{n} \log Z_n(\Phi, r) \leq 0.$$

Therefore, $r \geq s(\Phi)$ and the desired inequality follows from (5.15).

It remains to verify that $s(\Phi_*) \leq s(\Phi)$. Recall Lemma 2.4(i) which says that if $\Psi \in \Xi_X(K, \gamma, \theta)$ then for every $t > 0$ and every $n \geq 1$,

$$\mathrm{P}_\Psi(t) \leq \frac{1}{n} \log Z_n(\Psi, t) \leq \mathrm{P}_\Psi(t) + \frac{t \log K}{n}. \tag{5.17}$$

It is enough to prove that if $r > s(\Phi)$ then $r \geq s(\Phi_*)$. We have $P_\Phi(r) = 0$. Fix an arbitrary $\delta > 0$. It suffices to show that $P_{\Phi_*}(r) \leq \delta$. It is convenient to set $\Psi := \Phi_*$. By the first part of this proposition, $\Phi_* \in \Xi_X(K_1, \gamma, \theta)$. Take $q \in \mathbb{N}$ so large that $\frac{r \log K_1}{q} < \frac{\delta}{8}$. Then there exists a finite subset $F \subset I$ such that

$$0 < \frac{1}{q} \log Z_q(\Psi, r) - \frac{1}{q} \log Z_q(\Psi_F, r) < \frac{\delta}{4}.$$

Now applying (5.17) to Ψ and Ψ_F (which is obviously in $\Xi_X(K_1,\gamma,\theta)$ as well) we obtain for all $n \geq q$:

$$0 < \frac{1}{n}\log Z_n(\Psi,r) - \frac{1}{n}\log Z_n(\Psi_F,r) < \frac{\delta}{4} + \frac{2r\log K_1}{q} < \frac{\delta}{2}\,. \tag{5.18}$$

Recall that Ψ_F is a finite IFS so its elements (which are of the form $\phi_k^l\phi_i$) have some finite maximal length, say L, over the alphabet $\mathcal{A}$. Therefore, $Z_n(\Psi_F,r) \leq \sum_{p\leq nL} Z_p(\Phi,r)$. Since $P_\Phi(r)=0$, there exists $m\in\mathbb{N}$ such that $Z_p(\Phi,r)\leq e^{\frac{p\delta}{2L}}$ for $p\geq m$. Then

$$Z_n(\Psi_F,r) \leq \sum_{p\leq m} Z_p(\Phi,r) + \sum_{p=m+1}^{nL} e^{\frac{p\delta}{2L}} \leq C(e^{\delta n/2}).$$

This implies that $\lim_{n\to\infty}\frac{1}{n}\log Z_n(\Psi_F,r)\leq\frac{\delta}{2}$, and together with (5.18) this yields

$$\mathrm{P}_\Psi(r) = \lim_{n\to\infty}\frac{1}{n}\log Z_n(\Psi,r) \leq \frac{\delta}{2}+\frac{\delta}{2} = \delta,$$

as desired. The proof is complete. □

6. Families of parabolic IFS.

Let $U\subset\mathbb{R}^d$ be an open set. Here we consider families of parabolic IFS

$$\Phi^{\mathbf{t}} = \{\phi_1^{\mathbf{t}},\ldots,\phi_{k-1}^{\mathbf{t}},\phi_k\}$$

depending on $\mathbf{t}\in\overline{U}$, the closure of U. We assume that $\Phi^{\mathbf{t}}\in\Gamma_X(\theta)$ for all $\mathbf{t}\in\overline{U}$ (see Definition 5.5.). Although the parabolic function *does not* depend on the parameter, it is sometimes convenient to write $\phi_k^{\mathbf{t}}\equiv\phi_k$ for $\mathbf{t}\in U$. Let $\pi_{\mathbf{t}}:\ \mathcal{A}^\infty\to\mathbb{R}$ be the natural projection associated with $\Phi^{\mathbf{t}}$ and denote $J_{\mathbf{t}}=\pi_{\mathbf{t}}(\mathcal{A}^\infty)$. Two conditions which control the dependence on $\mathbf{t}$ will be needed.

Continuity: the maps

$$\mathbf{t}\mapsto\phi_i^{\mathbf{t}} \ \text{ are continuous from } \ \overline{U} \ \text{ to } \ \mathcal{C}^{1+\theta}(X) \ \text{ for } \ i\leq k-1. \tag{6.1}$$

Transversality Condition: there exists a constant C_1' such that for all ω and τ in $\mathcal{A}^\infty$ with $\omega_1\neq\tau_1$,

$$\mathcal{L}_d\{\mathbf{t}\in U:\ |\pi_{\mathbf{t}}(\omega)-\pi_{\mathbf{t}}(\tau)|\leq r\}\leq C_1' r \quad \text{for all } \ r>0. \tag{6.2}$$

This condition is almost identical to the transversality condition (3.2) except that here $\mathcal{A}$ is finite so C_1' is an absolute constant. Let

$$s(\mathbf{t}) = s(\Phi^{\mathbf{t}}) = \min\{u>0:\ P_{\Phi^{\mathbf{t}}}(u)=0\}\,.$$

The following theorem is the main result of the paper.

Theorem 6.1. *Suppose that* $\{\Phi^{\mathbf{t}}\}_{\mathbf{t}\in\overline{U}}$ *is a family of parabolic IFS in* $\Gamma_X(\theta)$ *satisfying (6.1) and (6.2). Then the function* $\mathbf{t}\mapsto s(\mathbf{t})$ *is continuous on* U *and*

(i) $\dim_{\mathrm{H}}(J_{\mathbf{t}}) = \min\{s(\mathbf{t}), 1\}$ *for Lebesgue-a.e.* $\mathbf{t}\in U$*;*

(ii) $\mathcal{L}_1(J_{\mathbf{t}}) > 0$ *for Lebesgue-a.e.* $\mathbf{t}\in U$ *such that* $s(\mathbf{t}) > 1$.

Proof. The plan is to apply Theorem 3.1 to the family of associated infinite hyperbolic IFS

$$\Phi^{\mathbf{t}}_* = \{\phi^n_k\phi^{\mathbf{t}}_i :\ n\geq 0,\ i\leq k-1\}.$$

Since the family $\{\Phi^{\mathbf{t}}\}_{\mathbf{t}\in\overline{U}}$ satisfies the continuity condition (6.1), we can find V,γ,u,M, so that

$$\Phi^{\mathbf{t}}\in\Gamma_X(\theta,V,\gamma,u,M)\quad\text{ for all }\ \mathbf{t}\in U, \tag{6.3}$$

see (5.4) for the meaning of this notation. These numbers and the neighborhood V of v will be fixed for the rest of the proof.

Observe that the constant K_1 in Corollary 5.9 is independent of $\mathbf{t}\in U$ and by Corollary 5.11(i) we have $\Phi^{\mathbf{t}}_*\in\Xi_X(K_1,\gamma,\theta)$ for $\mathbf{t}\in U$. Notice also that $s(\mathbf{t}) = s(\Phi^{\mathbf{t}}_*)$ by Proposition 5.10(ii). Thus, to prove the theorem it remains to verify the distortion continuity property (3.1) and transversality condition (3.2) for $\Phi^{\mathbf{t}}_*$. We begin with the latter since it is easier.

As in Section 5, the alphabet for the IFS $\Phi^{\mathbf{t}}_*$ will be $I = \{(n,i) :\ n\geq 0,\, i\leq k-1\}$. Let ζ and ξ be elements of I^∞ with $\zeta_1\neq\xi_1$. To distinguish between the IFS $\Phi^{\mathbf{t}}$ and $\Phi^{\mathbf{t}}_*$ we denote by $\pi^*_{\mathbf{t}} :\ I^\infty\to\mathbb{R}$ the natural projection corresponding to the IFS $\Phi^{\mathbf{t}}_*$. Let $\zeta_1 = (n,i)$ and $\xi_1 = (m,j)$. Clearly,

$$\pi^*_{\mathbf{t}}(\zeta) = \pi_{\mathbf{t}}(\omega),\quad \pi^*_{\mathbf{t}}(\xi) = \pi_{\mathbf{t}}(\tau)$$

where $\omega\in\mathcal{A}^\infty$ begins with $k^n i$ and $\tau\in\mathcal{A}^\infty$ begins with $k^m j$. Assume without loss of generality that $n\leq m$. Then

$$\pi^*_{\mathbf{t}}(\zeta) - \pi^*_{\mathbf{t}}(\xi) = \pi_{\mathbf{t}}(\omega) - \pi_{\mathbf{t}}(\tau) = \phi^n_k\left(\pi_{\mathbf{t}}(\sigma^n\omega)\right) - \phi^n_k\left(\pi_{\mathbf{t}}(\sigma^n\tau)\right)$$

and $(\sigma^n\omega)_1\neq(\sigma^n\tau)_1$. Thus, by the Mean Value Theorem and (5.7), for some $c\in X$,

$$|\pi^*_{\mathbf{t}}(\zeta) - \pi^*_{\mathbf{t}}(\xi)| = |(\phi^n_k)'(c)|\cdot|\pi_{\mathbf{t}}(\sigma^n\omega) - \pi_{\mathbf{t}}(\sigma^n\tau)|\geq u^n|\pi_{\mathbf{t}}(\sigma^n\omega) - \pi_{\mathbf{t}}(\sigma^n\tau)|\,.$$

Therefore, by (6.2),

$$\mathcal{L}_d\{\mathbf{t}\in U :\ |\pi^*_{\mathbf{t}}(\zeta) - \pi^*_{\mathbf{t}}(\xi)|\leq r\}\leq\ \mathcal{L}_d\{\mathbf{t}\in U :\ |\pi_{\mathbf{t}}(\sigma^n\omega) - \pi_{\mathbf{t}}(\sigma^n\tau)|\leq ru^{-n}\}\leq C_1' ru^{-n}$$

which implies (3.2) with $C(\xi_1,\eta_1) := C_1' u^{-n}$.

The property (3.1) for $\Phi_*^{\mathtt{t}}$ is immediate from (6.1) and the following lemma where we let

$$\| \Phi - \Psi \| = \max_{i \le k-1} \| \phi_i - \psi_i \| \quad \text{and} \quad \| \Phi' - \Psi' \| = \max_{i \le k-1} \| \phi_i' - \psi_i' \| .$$

Lemma 6.2. *There exists a positive constant $C_2' = C_2'(X, \theta, V, \gamma, u, M)$ such that for any $\Phi = \{\phi_1, \dots, \phi_k\}$ and $\Psi = \{\psi_1, \dots, \psi_k\}$, two parabolic IFS in $\Gamma_X(\theta, V, \gamma, u, M)$ with $\phi_k = \psi_k$, the associated infinite hyperbolic IFS Φ_* and Ψ_* have the following property. For any $\tau \in I^*$ and any $x \in X$,*

$$\frac{|(\phi_\tau^*)'(x)|}{|(\psi_\tau^*)'(x)|} \le \exp\left(C_2' |\tau| \left(\| \Phi - \Psi \|^\theta + \| \Phi' - \Psi' \| \right) \right) .$$

Proof. Let $n = |\tau|$ and observe that for all $x \in X$ and $0 \le j < n$:

$$\begin{aligned} &|\phi^*_{\sigma^j\tau}(x) - \psi^*_{\sigma^j\tau}(x)| \\ &\quad \le |\phi^*_{\sigma^j(\tau|_{n-1})}(\phi^*_{\tau_n}(x)) - \phi^*_{\sigma^j(\tau|_{n-1})}(\psi^*_{\tau_n}(x))| + |\phi^*_{\sigma^j(\tau|_{n-1})}(\psi^*_{\tau_n}(x)) - \psi^*_{\sigma^j(\tau|_{n-1})}(\psi^*_{\tau_n}(x))| \\ &\quad \le \gamma^{n-j-1} \| \Phi_* - \Psi_* \| + |\phi^*_{\sigma^j(\tau|_{n-1})}(x') - \psi^*_{\sigma^j(\tau|_{n-1})}(x')|, \end{aligned}$$

for $x' = \psi^*_{\tau_n}(x)$. Proceeding inductively we obtain

$$|\phi^*_{\sigma^j\tau}(x) - \psi^*_{\sigma^j\tau}(x)| \le \sum_{i=1}^{n-j} \gamma^{n-i-1} \| \Phi_* - \Psi_* \| < \frac{\| \Phi_* - \Psi_* \|}{1-\gamma} . \tag{6.4}$$

Observe that

$$\| \Phi_* - \Psi_* \| = \max_{i \le k-1,\ m \ge 0} \| \phi_k^m \phi_i - \phi_k^m \psi_i \| = \max_{i \le k-1} \| \phi_i - \psi_i \| = \| \Phi - \Psi \| \tag{6.5}$$

since $\| (\phi_k^m)' \| = 1$. Now let $y = \phi^*_{\sigma^j\tau}(x)$ and $z = \psi^*_{\sigma^j\tau}(x)$, and suppose that $\tau_j = (l, i)$. Then $\phi^*_{\tau_j} = \phi_k^l \phi_i$ and $\psi^*_{\tau_j} = \phi_k^l \psi_i$, and we can estimate by (5.9):

$$\begin{aligned} \log \left| \frac{(\psi^*_{\tau_j})'(\psi^*_{\sigma^j\tau}(x))}{\left(\phi^*_{\tau_j}\right)'(\phi^*_{\sigma^j\tau}(x))} \right| &= \log \left| \frac{(\phi_k^l \psi_i)'(z)}{(\phi_k^l \phi_i)'(y)} \right| = \log \left| \frac{(\phi_k^l)'(\psi_i(z))}{(\phi_k^l)'(\phi_i(y))} \right| + \log \left| \frac{\psi_i'(z)}{\phi_i'(y)} \right| \\ &\le \log L_5 \cdot |\psi_i(z) - \phi_i(y)|^\theta + u^{-1} |\psi_i'(z) - \phi_i'(y)|. \end{aligned} \tag{6.6}$$

Next,

$$\begin{aligned} |\psi_i(z) - \phi_i(y)| &\le |\psi_i(z) - \psi_i(y)| + |\psi_i(y) - \phi_i(y)| \\ &\le |z - y| + \| \Phi - \Psi \| \le \| \Phi - \Psi \| \ \frac{2-\gamma}{1-\gamma} \end{aligned}$$

by (6.4) and (6.5). Furthermore,

$$\begin{aligned} |\psi_i'(z) - \phi_i'(y)| &\le |\psi_i'(z) - \psi_i'(y)| + |\psi_i'(y) - \phi_i'(y)| \\ &\le \| \Psi' \|_\theta \ |z - y|^\theta + \| \Phi' - \Psi' \| \\ &\le M \| \Phi - \Psi \|^\theta \ (1-\gamma)^{-\theta} + \| \Phi' - \Psi' \| . \end{aligned}$$

Combining the last two estimates with (6.6) yields

$$\begin{aligned}
&\log\left|\frac{(\psi^*_{\tau_j})'(\psi^*_{\sigma^j\tau}(x))}{\left(\phi^*_{\tau_j}\right)'(\phi^*_{\sigma^j\tau}(x))}\right| \\
&\leq \| \Phi - \Psi \|^\theta \left(\frac{2-\gamma}{1-\gamma}\right)^\theta \log L_5 + u^{-1}\left(M \| \Phi - \Psi \|^\theta (1-\gamma)^{-\theta} + \| \Phi' - \Psi' \|\right) \\
&\leq C_2'\left(\| \Phi - \Psi \|^\theta + \| \Phi' - \Psi' \|\right)
\end{aligned}$$

for some constant $C_2' = C_2'(X, \theta, V, \gamma, u, M)$. Exchanging the roles of Φ and Ψ we can see that

$$\left|\log\left|\frac{(\psi^*_{\tau_j})'(\psi^*_{\sigma^j\tau}(x))}{\left(\phi^*_{\tau_j}\right)'(\phi^*_{\sigma^j\tau}(x))}\right|\right| \leq C_2'\left(\| \Phi - \Psi \|^\theta + \| \Phi' - \Psi' \|\right).$$

Finally,

$$\begin{aligned}
\left|\log\left|\frac{(\psi^*_{\tau})'(x))}{(\phi^*_{\tau})'(x))}\right|\right| &= \left|\sum_{j=1}^n \log\left|\frac{(\psi^*_{\tau_j})'(\psi^*_{\sigma^j\tau}(x))}{(\phi^*_{\tau_j})'(\phi^*_{\sigma^j\tau}(x))}\right|\right| \\
&\leq \sum_{j=1}^n C_2'\left(\| \Phi - \Psi \|^\theta + \| \Phi' - \Psi' \|\right) \\
&= nC_2'\left(\| \Phi - \Psi \|^\theta + \| \Phi' - \Psi' \|\right).
\end{aligned}$$

The lemma is proved, and this also concludes the proof of Theorem 6.1. □

Given $\omega \in \mathcal{A}^*$ let $h(\omega)$ denote the number of hyperbolic letters (i.e.$\neq k$) appearing in ω. We record the following useful corollary for future reference although it is not needed in this paper.

Corollary 6.3. *There exists a positive constant $C_2' = C_2'(X, \theta, V, \gamma, u, M)$ such that for any $\Phi = \{\phi_1, \ldots, \phi_k\}$ and $\Psi = \{\psi_1, \ldots, \psi_k\}$, two parabolic IFS in $\Gamma_X(\theta, V, \gamma, u, M)$ with $\phi_k = \psi_k$, for all $\omega \in \mathcal{A}^*$ and all $x \in X$,*

$$\frac{|\phi'_\omega(x)|}{|\psi'_\omega(x)|} \leq \exp\left(C_2' h(\omega)\left(\| \Phi - \Psi \|^\theta + \| \Phi' - \Psi' \|\right)\right).$$

Proof. For any $\omega \in \mathcal{A}^*$ we can find $\tau \in I^*$ and $l \geq 0$ such that

$$\phi_\omega = \phi^*_\tau \phi^l_k \qquad \text{and} \qquad \psi_\omega = \psi^*_\tau \phi^l_k.$$

By Lemma 6.2 we have

$$\begin{aligned}
\frac{|\phi'_\omega(x)|}{|\psi'_\omega(x)|} = \frac{|(\phi^*_\tau)'(\phi^l_k(x))(\phi^l_k)'(x)|}{|(\psi^*_\tau)'(\phi^l_k(x))(\phi^l_k)'(x)|} &= \frac{|(\phi^*_\tau)'(\phi^l_k(x))|}{|(\psi^*_\tau)'(\phi^l_k(x))|} \\
&\leq \exp\left(C_2'|\tau|\left(\| \Phi - \Psi \|^\theta + \| \Phi' - \Psi' \|\right)\right)
\end{aligned}$$

which finishes the proof since $|\tau| = h(\omega)$. □

We conclude this section with some finer results concerning exceptional parameters in Theorem 6.1(i). These will turn out to be almost immediate consequences of the results obtained in Section 4. The strong transversality condition is formulated in the context of parabolic systems in the same manner as the strong transversality condition in Section 4. Recall that $N_r(F)$ is the minimal number of balls of radius r needed to cover the set $F \subset \mathbb{R}^d$.

STRONG TRANSVERSALITY CONDITION: there exists a constant C_1' such that for all ω and τ in $\mathcal{A}^\infty$ with $\omega_1 \neq \tau_1$,

$$N_r\left(\{\mathbf{t} \in U : \ |\pi_{\mathbf{t}}(\omega) - \pi_{\mathbf{t}}(\tau)| \le r\}\right) \le C_1' r \quad \text{for all} \ \ r > 0. \tag{6.7}$$

The same arguments as used in the proof of Theorem 6.1 demonstrate that the strong transversality condition (6.7) for a parabolic IFS implies the strong transversality condition (4.1) for the associated hyperbolic system. Consequently, Theorem 4.3 and Corollary 4.4 respectively imply the following.

Theorem 6.4. *Suppose that $\{\Phi^{\mathbf{t}}\}_{\mathbf{t}\in\overline{U}}$ is a d-parameter family of parabolic IFS in $\Gamma_X(\theta)$ satisfying (6.1) and (6.7). If G is an arbitrary subset of U, then for every $\xi > 0$ we have*

$$\dim_{\mathrm{H}}\left(\left\{\mathbf{t} \in G : \ \dim_{\mathrm{H}}(J_t) < \min\{\xi, s(\mathbf{t})\}\right\}\right) \le \min\{\xi, \sup_G s(\mathbf{t})\} + d - 1.$$

Corollary 6.5. *Suppose that $\{\Phi^{\mathbf{t}}\}_{\mathbf{t}\in\overline{U}}$ is a d-parameter family of parabolic IFS in $\Gamma_X(\theta)$ satisfying (6.1) and (6.7). Then for every $\mathbf{t}_0 \in U$ we have*

$$\lim_{r\to 0} \dim_{\mathrm{H}}\left(\{t \in B(\mathbf{t}_0, r) : \dim_{\mathrm{H}}(J_{\mathbf{t}}) < \min\{\xi, s(\mathbf{t})\}\}\right) \le \min\{\xi, s(\mathbf{t}_0)\} + d - 1.$$

7. EXAMPLES

7.1. **General classes of examples.** We are going to apply Theorem 6.1 to two more specific types of IFS families. In each case we need to impose some bounds on derivatives to guarantee transversality (even strong transversality), which is the most difficult condition to check. Recall that $\|\cdot\|_Y$ denotes the supremum norm on Y.

Proposition 7.1. *Let $\Phi = \{\phi_1, \ldots, \phi_k\} \in \Gamma_X(\theta)$ and $1 \le d \le k-1$. Assume that*

$$\phi_i \ \ \text{are increasing for all} \ \ i \le k; \tag{7.1}$$

$$\phi_i(X) \cap \phi_j(X) = \emptyset \ \ \text{for all} \ \ d < i < j \le k. \tag{7.2}$$

and

$$\| \phi_i' \| + \| \phi_j' \|_{\phi_j^{-1}\phi_i(X)} < 1 \quad \textit{for all} \quad i < j \quad \textit{such that} \quad \phi_i(X) \cap \phi_j(X) \neq \emptyset. \tag{7.3}$$

Consider the family

$$\Phi^{\mathbf{t}} = \{\phi_1(x) + t_1, \ldots, \phi_d(x) + t_d, \phi_{d+1}(x), \ldots, \phi_k(x)\} \quad \textit{where} \quad \mathbf{t} = (t_1, \ldots, t_d) \in \mathbb{R}^d.$$

Then there exists $\eta > 0$ such that $\{\Phi^{\mathbf{t}} : \mathbf{t} \in B_\eta(\mathbf{0})\}$ satisfies all the hypotheses of Theorem 6.1. Moreover, the strong transversality condition (6.7) is satisfied.

Remark. It is not hard to find specific IFS satisfying the above proposition. Notice that non-trivial examples start with $k = 3$. Indeed, if $k = 2$ there is a dichotomy: the limit set is either connected or the projection map $\pi_\Phi : \mathcal{A}^\infty \to J_\Phi$ is one-to-one. In the latter case there are no "overlaps", so the dimension formula holds by [U]. In the former case J_Φ is an interval, so there is no question about its dimension or measure. To get a non-trivial example with $k = 3$, one can take a parabolic function, say, $\phi_3(x) = \sin x$ on $[0, b]$, with $b < \pi/2$, and add two more functions so that there is a "gap" and an "overlap". For instance, an increasing function ϕ_2 may be chosen to satisfy $\phi_2(0) > \sin b$ and $\phi_2(b) = b$, and the increasing function ϕ_1 to satisfy $\phi_1(0) \in (0, \sin b)$, $\phi_1(b) < \phi_2(0)$. Consider a one-parameter family $\Phi^t = \{\phi_1(x) + t, \phi_2(x), \phi_3(x)\}$ and call t admissible if these properties persist for Φ^t. If t is admissible, then the limit set has the convex hull equal to $[0, b]$. We can define $X = [0, b + \varepsilon]$ for a small $\varepsilon > 0$, so that $\phi_2(X) \subset \mathrm{Int}(X)$. All the assumptions of Proposition 7.1 are satisfied, provided that the derivative of ϕ_1 is sufficiently small, so the results of Section 6 apply.

Example 1. The family of IFS

$$\Phi^t = \{0.01(x+1)^2 + t,\ 0.1x + (4.5/11)\pi,\ \sin x\} \quad \text{on} \quad [0, \tfrac{5\pi}{11} + \varepsilon]$$

satisfies all the conditions above for $t \in [0.3, 0.97]$ if $\varepsilon >$ is sufficiently small. (We are grateful to P. Hanus for correcting our original faulty example.)

In the next proposition all elements of the IFS are assumed to be of the form $\phi(x + a_j)$ for a single function ϕ. Although this seems much more special than Proposition 7.1, it covers some interesting families, in particular, those which arise from a class of continued fractions.

Proposition 7.2. (i) *Let $\Phi = \{\phi_1, \ldots, \phi_k\} \in \Gamma_X(\theta)$ and $1 \leq d \leq k-1$ be such that (7.2) is satisfied. We further assume that there exists a single increasing function $\phi \in \mathcal{C}^{1+\theta}$ on*

some interval Y *and* $a_i \in \mathbb{R}$, $u > 0$ *such that*

$$\phi_i(x) = \phi(x + a_i), \ i \le k, \tag{7.4}$$

and

$$\inf_{x \in Y} |\phi'(x)| \ge u > 0. \tag{7.5}$$

Let

$$\Phi^{\mathbf{t}} = \{\phi_1(x + t_1), \dots, \phi_d(x + t_d), \phi_{d+1}(x), \dots, \phi_k(x)\} \quad \textit{where} \ \ \mathbf{t} = (t_1, \dots, t_d) \in \mathbb{R}^d.$$

Denote by $\pi_{\mathbf{t}}$ *the projection map corresponding to* $\Phi^{\mathbf{t}}$. *If there exists* $\delta > 0$ *such that*

$$\frac{\partial}{\partial t_i} \pi_{\mathbf{t}}(\xi)|_{\mathbf{t}=0} < 1 - \delta \quad \textit{for all} \ \ \xi \in \mathcal{A}^\infty \ \ \textit{and all} \ \ 1 \le i \le d, \tag{7.6}$$

then there exists $\eta > 0$ *such that* $\{\Phi^{\mathbf{t}} : \ \mathbf{t} \in B_\eta(\mathbf{0})\}$ *satisfies all the hypotheses of Theorem 6.1. Moreover, the strong transversality condition (6.7) is satisfied.*

(ii) *Condition (7.6) holds if*

$$\| \phi_i' \| < 1/2 \quad \textit{for all} \ \ 1 \le i \le d. \tag{7.7}$$

Example 2. Let $\Phi^a = \{\ln(x + a), \ \ln(x + e^2 - 2), \ \ln(x + 1)\}$. This is a parabolic IFS on $[0, 2]$. We have $\| \phi_1' \| = 1/a$, and so Proposition 7.2 and Theorem 6.1 imply that the dimension formula holds for a.e. $a \in (2, e^2 - 2)$. (We need to enlarge the interval slightly and let $X = [0, 2 + \varepsilon]$ so that $\phi_2(X) \subset \mathrm{Int}(X)$.)

Turning to the proofs of Propositions 7.1 and 7.2, note that checking strong transversality (6.7) is the only issue since all other properties obviously hold for sufficiently small perturbations of Φ. This will be done with the help of the following elementary lemma.

Lemma 7.3. *Let* $U \subset \mathbb{R}^d$ *be an open, bounded set with smooth boundary. Suppose that* f *is a* $\mathcal{C}^1$ *real-valued function in a neighborhood of* $\overline{U}$ *such that for some* $i \in \{1, \dots, d\}$ *there exists* $\eta > 0$ *satisfying*

$$\mathbf{t} \in U, \quad |f(\mathbf{t})| \le \eta \implies \frac{\partial f(\mathbf{t})}{\partial t_i} \ge \eta. \tag{7.8}$$

Then there exists $C = C(\eta)$ *such that*

$$N_r(\{\mathbf{t} \in U : \ |f(\mathbf{t})| \le r\}) \le Cr, \quad \forall \ r > 0. \tag{7.9}$$

Proof. Recall that $N_r(F)$ denotes the minimal number of balls of radius r needed to cover F. Since U is bounded it suffices to prove our lemma for every $r < \eta$. Let $M = f^{-1}(0)$. By (7.8), $\mathrm{grad}(f)(\mathbf{t}) \ne 0$ for every $\mathbf{t} \in M$. Hence M is a $(d-1)$-dimensional $\mathcal{C}^1$-manifold and $(M \cap \overline{U}) \cup \partial U$ is contained in a union of finitely many compact connected $\mathcal{C}^1$ manifolds

with smooth boundaries. Thus there exists a constant C_1 such that for every $r < \eta$ the set $(M \cap \overline{U}) \cup \partial U$ can be covered by $C_1 r^{1-d}$ balls with radii r. In order to complete the proof it is sufficient to show that for each point $\mathbf{t} = (t_1, t_2, \dots, t_d) \in U$ with $|f(\mathbf{t})| \leq r$, the distance between $\mathbf{t}$ and $(M \cap \overline{U}) \cup \partial U$ does not exceed r/η. Indeed, we will then cover the set $f^{-1}([-r, r])$ by at most $C_1 r^{1-d}$ balls with radii $(1 + 1/\eta)r$. But each ball of radius $(1 + 1/\eta)r$ can be covered by $C_2(\eta)$ balls of radii r, where $C_2(\eta)$ depends only on η and the dimension d and we would therefore be done.

Without loss of generality let $i = 1$. Consider the function $g(s) = f(s, t_2, \dots, t_d)$ defined in a neighborhood of t_1. Suppose first that $g(t_1) \in [-r, 0]$. By (7.8), the function $g(s)$ increases for $s > t_1$ until either the point $(s, t_2, \dots, t_d)$ reaches the boundary of U, say at a point $(u, t_2, \dots, t_d) \in \partial U$ and $g(s) \in [-r, 0]$ for all $s \in [t_1, u]$, or $g(s)$ will take on the value zero earlier. Suppose that the first case materializes. Then $g'(s) \geq \eta$ for every $s \in [t_1, u]$ by (7.8). By the Mean Value Theorem, $|u - t_1| \leq r/\eta$, and therefore, $\operatorname{dist}(\mathbf{t}, \partial U) \leq \operatorname{dist}(\mathbf{t}, (u, t_2, \dots, t_d)) = |u - t_1| \leq r/\eta$. So, we are done in this case. If the second case holds, let w be that point for which $g(w) = 0$ and $g(s) \in [-r, 0]$ for every $s \in [t_1, w]$. Then $(w, t_2, \dots, t_d) \in M$ and $g'(s) \geq \eta$ for every $s \in [t_1, w]$. Again by the Mean Value Theorem, $|w - t_1| \leq r/\eta$. Then $\operatorname{dist}(\mathbf{t}, M \cap \overline{U}) \leq \operatorname{dist}(\mathbf{t}, (w, t_2, \dots, t_d)) = |w - t_1| \leq r/\eta$ and we are done in this case as well. If $g(t_1) \in [0, r]$, we proceed similarly letting s go left from t_1. The proof is complete. □

Proof of Proposition 7.1. It is convenient to write $\Phi^{\mathbf{t}} = \{\phi_i^{\mathbf{t}}\}_{i \leq k}$. The properties (7.3) and (5.1) persist under a small perturbation, so we can find $\eta \in (0, 1)$ such that

$$\inf\{\operatorname{dist}(\phi_i^{\mathbf{t}}(X), v) : \ i \leq k - 1, \ \| \mathbf{t} \| \leq \eta\} \geq \eta > 0; \tag{7.10}$$

$$\inf\{\operatorname{dist}\left(\phi_i^{\mathbf{t}}(X), \phi_j^{\mathbf{t}}(X)\right) : \ \phi_i(X) \cap \phi_j(X) = \emptyset, \ \| \mathbf{t} \| \leq \eta\} \geq \eta > 0; \tag{7.11}$$

$$\operatorname{dist}(\phi_i(X), \phi_j(x)) \leq 3\eta \implies \| \phi_i' \| + |\phi_j'(x)| < 1 - \eta \quad \text{for} \ \ i \leq d, \ i < j. \tag{7.12}$$

The continuity property (6.1) is obvious so we only need to verify the strong transversality condition (6.7).

Consider $\omega, \tau \in \mathcal{A}^\infty$, with $\omega_1 \neq \tau_1$. Let $i = \omega_1$ and $j = \tau_1$. Without loss of generality assume that $i < j$. If $\phi_i(X) \cap \phi_j(X) = \emptyset$ then

$$|\pi_{\mathbf{t}}(\omega) - \pi_{\mathbf{t}}(\tau)| = |\phi_i^{\mathbf{t}}(\pi_{\mathbf{t}}(\sigma\omega)) - \phi_j^{\mathbf{t}}(\pi_{\mathbf{t}}(\sigma\tau))| \geq \eta > 0$$

for $\mathbf{t} \in B_\eta(\mathbf{0})$ by (7.11) and

$$\{\mathbf{t} \in B_\eta(\mathbf{0}) : \ |\pi_{\mathbf{t}}(\omega) - \pi_{\mathbf{t}}(\tau)| \le r\} \subset B_r(\mathbf{0}) \tag{7.13}$$

(note that the left-hand side is empty for $r < \eta$). Thus (6.7) certainly holds in this case. If $\phi_i(X) \cap \phi_j(X) \neq \emptyset$ then $i \le d$ and

$$\pi_{\mathbf{t}}(\omega) - \pi_{\mathbf{t}}(\tau) = t_i + \phi_i(\pi_{\mathbf{t}}(\sigma\omega)) - \phi_j^{\mathbf{t}}(\pi_{\mathbf{t}}(\sigma\tau)).$$

We are going to use Lemma 7.3 with $f(\mathbf{t}) = \pi_{\mathbf{t}}(\omega) - \pi_{\mathbf{t}}(\tau)$ and $U = B_\eta(\mathbf{0})$, so we need to check (7.8). Since $|\phi_j^{\mathbf{t}}(\pi_{\mathbf{t}}(\sigma\tau)) - \phi_j(\pi_{\mathbf{t}}(\sigma\tau))| = |t_j| \le \eta$ and $|t_i| \le \eta$, we have the implication

$$\begin{aligned} |f(\mathbf{t})| = |\pi_{\mathbf{t}}(\omega) - \pi_{\mathbf{t}}(\tau)| \le \eta \ &\implies \ |\phi_i(\pi_{\mathbf{t}}(\sigma\omega)) - \phi_j(\pi_{\mathbf{t}}(\sigma\tau))| \le 3\eta \\ &\implies \| \phi_i' \| + |\phi_j'(\pi_{\mathbf{t}}(\sigma\tau))| < 1 - \eta \end{aligned} \tag{7.14}$$

by (7.12). Since $j \neq i$ we have

$$\begin{aligned} \frac{\partial}{\partial t_i}(\pi_{\mathbf{t}}(\omega) - \pi_{\mathbf{t}}(\tau)) &= 1 + \frac{\partial}{\partial t_i}\phi_i(\pi_{\mathbf{t}}(\sigma\omega)) - \frac{\partial}{\partial t_i}\phi_j(\pi_{\mathbf{t}}(\sigma\tau)) \\ &= 1 + \phi_i'(\pi_{\mathbf{t}}(\sigma\omega)) \cdot \frac{\partial}{\partial t_i}\pi_{\mathbf{t}}(\sigma\omega) - \phi_j'(\pi_{\mathbf{t}}(\sigma\tau)) \cdot \frac{\partial}{\partial t_i}\pi_{\mathbf{t}}(\sigma\tau). \end{aligned} \tag{7.15}$$

Observe that $\phi_i' > 0$ and $\frac{\partial}{\partial t_i}\pi_{\mathbf{t}}(\sigma\omega) \ge 0$ by the assumption (7.1). Suppose that $(\sigma\tau)_n$ is the first symbol i in $\sigma\tau$ (if $\sigma\tau$ contains no i we have $\frac{\partial}{\partial t_i}\pi_{\mathbf{t}}(\sigma\tau) = 0$ and the claim is obvious). Then

$$\pi_{\mathbf{t}}(\sigma\tau) = \phi^{\mathbf{t}}_{\sigma\tau|_{n-1}}\left(t_i + \phi_i\left(\pi_{\mathbf{t}}(\sigma^{n+1}\tau)\right)\right),$$

and since $\| (\phi^{\mathbf{t}}_{\sigma\tau|_{n-1}})' \| \le 1$ we obtain

$$\frac{\partial}{\partial t_i}\pi_{\mathbf{t}}(\sigma\tau) \le 1 + \| \phi_i' \| \cdot \frac{\partial}{\partial t_i}\pi_{\mathbf{t}}(\sigma^{n+1}\tau).$$

Proceeding inductively we see that

$$\frac{\partial}{\partial t_i}\pi_{\mathbf{t}}(\sigma\tau) \le 1 + \| \phi_i' \| + \| \phi_i' \|^2 + \ldots = (1 - \| \phi_i' \|)^{-1}$$

(recall that ϕ_i is hyperbolic since $i \le d < k$). Together with (7.15) this implies

$$\begin{aligned} \frac{\partial}{\partial t_i}(\pi_{\mathbf{t}}(\omega) - \pi_{\mathbf{t}}(\tau)) &\ge 1 - |\phi_j'(\pi_{\mathbf{t}}(\sigma\tau))|(1 - \| \phi_i' \|)^{-1} \\ &= \frac{1 - \| \phi_i' \| - |\phi_j'(\pi_{\mathbf{t}}(\sigma\tau))|}{1 - \| \phi_i' \|} > \eta \end{aligned} \tag{7.16}$$

by (7.14). We have verified (7.8), so the strong transversality condition (6.7) holds by Lemma 7.3 and the proof is complete. □

Proof of Proposition 7.2 is similar to that of Proposition 7.1. We can choose $\eta > 0$ so that (7.10) and (7.11) hold, and moreover,

$$\frac{\partial}{\partial t_i}\pi_{\mathbf{t}}(\xi) < 1-\eta \quad \text{for all } \ \xi \in \mathcal{A}^\infty,\ i \le d,\ \mathbf{t} \in B_\eta(\mathbf{0}). \tag{7.17}$$

Again we only need to check strong transversality. Let $\omega, \tau \in \mathcal{A}^\infty$ with $i = \omega_1, j = \tau_1$, and $i < j$. If $\phi_i(X) \cap \phi_j(X) = \emptyset$ we immediately obtain (7.13) by (7.11). If $\phi_i(X) \cap \phi_j(X) \neq \emptyset$ then $i \le d$ and

$$\begin{aligned}\pi_{\mathbf{t}}(\omega) - \pi_{\mathbf{t}}(\tau) &= \phi_i\left(t_i + \pi_{\mathbf{t}}(\sigma\omega)\right) - \phi_j^{\mathbf{t}}(\pi_{\mathbf{t}}(\sigma\tau)\Big) \\ &= \phi\left(a_i + t_i + \pi_{\mathbf{t}}(\sigma\omega)\right) - \phi\left(a_j + \kappa_j t_j + \pi_{\mathbf{t}}(\sigma\tau)\right) \\ &= \phi'(c)\left(A(t_j) + t_i + \pi_{\mathbf{t}}(\sigma\omega) - \pi_{\mathbf{t}}(\sigma\tau)\right)\end{aligned}$$

where $\kappa_j = 1$ if $j \le d$ and 0 otherwise, and $A(t_j) = a_i - a_j - \kappa_j t_j$ does not depend on t_i. Denoting $f(\mathbf{t}) := A(t_j) + t_i + \pi_{\mathbf{t}}(\sigma\omega) - \pi_{\mathbf{t}}(\sigma\tau)$ we have by (7.5):

$$\{\mathbf{t} \in B_\eta(\mathbf{0}) :\ |\pi_{\mathbf{t}}(\omega) - \pi_{\mathbf{t}}(\tau)| \le r\} \subset \{\mathbf{t} \in B_\eta(\mathbf{0}) :\ |f(\mathbf{t})| \le u^{-1}r\}\,.$$

If we show that $\frac{\partial f}{\partial t_i} \ge \eta$ then strong transversality will follow from Lemma 7.3. Since all ϕ_j are increasing it is easy to see that $\frac{\partial}{\partial t_i}\pi_{\mathbf{t}}(\sigma\omega) \ge 0$, hence

$$\frac{\partial f}{\partial t_i} \ge 1 - \frac{\partial}{\partial t_i}\pi_{\mathbf{t}}(\sigma\tau),$$

and the desired statement follows from (7.17).

(ii) It remains to derive (7.6) from (7.7). Let $1 \le i \le d$. If $\xi \in \mathcal{A}^\infty$ contains no i, then $\frac{\partial}{\partial t_i}\pi_{\mathbf{t}}(\xi) = 0$. Otherwise we write $\xi = wi\tilde{\xi}$ where $w \in \mathcal{A}^*$ contains no i (the word w may be empty). Then

$$\pi_{\mathbf{t}}(\xi) = \phi_w^{\mathbf{t}}\phi_i(t_i + \pi_{\mathbf{t}}(\tilde{\xi})),$$

hence

$$\frac{\partial}{\partial t_i}\pi_{\mathbf{t}}(\xi) \le \| \left(\phi_w^{\mathbf{t}}\right)' \| \cdot \| \phi_i' \| \left(1 + \frac{\partial}{\partial t_i}\pi_{\mathbf{t}}(\tilde{\xi})\right) \le \| \phi_i' \| \left(1 + \frac{\partial}{\partial t_i}\pi_{\mathbf{t}}(\tilde{\xi})\right).$$

Proceeding inductively we obtain that

$$\frac{\partial}{\partial t_i}\pi_{\mathbf{t}}(\tilde{\xi}) \le \sum_{n=1}^{\infty} \| \phi_i' \|^n = 1 - \frac{1 - 2\| \phi_i' \|}{1 - \| \phi_i' \|} < 1 - \delta$$

for some $\delta > 0$ by (7.7), and the proof is complete. □

7.2. **A class of continued fractions.** Here we study in detail IFS arising from the function $\phi(x) = \frac{x}{x+1}$. Let $A = \{a_1, \ldots, a_k\}$ where $a_i > 0$ for $i \le k-1$ and $a_k = 0$. Let $\Phi_A = \{\phi(x + a_i)\}_{i \le k}$. Then $\phi_k = \phi$ is parabolic on $[0,1]$, with the neutral fixed point $v = 0$, for all A. The functions $\phi(x+a_i)$, for $\imath \le k-1$, are hyperbolic on $[0,1]$. Clearly, the IFS Φ_A belongs to $\Gamma_{[0,1]}(1)$ (see Definition 5.5). The IFS Φ_A with $A = \{\alpha, 0\}$ was recently considered by Lyons [Ly] who studied the properties of the measure which arises from applying the maps randomly and independently with equal probabilities. However, as far as the limit set is concerned, the case $k = 2$ is trivial, since then either the limit set is an interval, or the IFS has no overlaps, so the dimension formula holds for all parameters by [U].

The connection with continued fractions is as follows: the limit set J_A of Φ_A on $[0,1]$ coincides with the set of continued fractions of the form

$$y = [1, Y_1, 1, Y_2, 1, Y_3, 1, \ldots] := \cfrac{1}{1 + \cfrac{1}{Y_1 + \cfrac{1}{1 + \ldots}}}$$

where $Y_i \in A$, and also with the set of continued fractions of the form

$$y = 1 - \cfrac{1}{(2 + Y_1) - \cfrac{1}{(2 + Y_2) - \cfrac{1}{(2 + Y_3) - \ldots}}}$$

where $Y_i \in A$. Indeed, these representations are immediate by writing

$$\frac{x+\alpha}{x+\alpha+1} = \cfrac{1}{1 + \cfrac{1}{\alpha + x}} \qquad \text{and} \qquad \frac{x+\alpha}{x+\alpha+1} = 1 - \frac{1}{(1+\alpha)+x}\,.$$

Example 3. Let $A = \{a, 2, 0\}$. We start with the one-parameter family of IFS Φ_A corresponding to $a \in (0,2)$, however, we will see that the nontrivial interval of parameters is smaller.

First observe that the limit set of the IFS $\{\phi(x+a),\ \phi(x)\}$ is an interval if $a \in (0, 1/2]$. Let $\phi_1(x) = \phi(x+a)$. Then $\phi_1(b) = b$ where $b = \frac{1}{2}(-a + \sqrt{a^2 + 4a})$. This implies that $[0,b] \supset \phi_1([0,b]) \cup \phi([0,b])$, and this becomes equality if and only if $\phi(b) \ge \phi_1(0)$ which is equivalent to $a \le 1/2$. Since the limit set can only increase when more functions are added, J_A contains an interval if $a \le 1/2$.

Since $\phi_2(x) = \phi(x+2)$ has the fixed point at $\sqrt{3}-1$, we see that $Y := [0, \sqrt{3}-1]$ is the convex hull of J_A. We have, denoting $\phi_3 = \phi$,

$$\phi_1(Y) = \left[\frac{a}{a+1}, \frac{\sqrt{3}-1+a}{\sqrt{3}+a}\right], \quad \phi_2(Y) = [2/3, \sqrt{3}-1], \quad \phi_3(Y) = [0, 1-1/\sqrt{3}].$$

Now, $\phi_2(Y) \cap \phi_3(Y) = \emptyset$. We want to make sure that the IFS has both a "gap" and an "overlap". Note that $\phi_1(Y) \cap \phi_2(Y) \neq \emptyset$ if and only if $a \in [3-\sqrt{3}, 2)$ and $\phi_1(Y) \cap \phi_3(Y) \neq \emptyset$ if and only if $a \in (0, \sqrt{3}-1]$ (recall that we only consider $0 < a < 2$). Thus the "interesting" set of parameters is $U := (\frac{1}{2}, \sqrt{3}-1) \cap (3-\sqrt{3}, 2)$.

We want to apply Proposition 7.2 with $k = 3$ and $d = 1$. Let $X = [0, \sqrt{3}-1+\varepsilon]$ where ε is so small that $\phi_2(X) \cap \phi_3(X) = \emptyset$. We have

$$\| \phi_1' \| = \phi_1'(0) = (1+a)^{-2} < 4/9 < 1/2,$$

since $a > \frac{1}{2}$. Thus (7.7) holds and all the assumptions of Proposition 7.2 are satisfied. Let $s(a) = s(\Phi_A)$ where $A = \{a, 2, 0\}$; then $s : U \to \mathbb{R}$ is continuous. Let $U_{<1} = \{A \in U : \ s(A) < 1\}$ and $U_{>1} = \{A \in U : \ s(A) > 1\}$. We obtain the following

Corollary 7.4. **(i)** $\dim_{\mathrm{H}}(J_A) = \min\{s(a), 1\}$ *for Lebesgue-a.e.* $a \in U$.

(ii) *For any set* $G \subset U_{<1}$ *we have*

$$\dim_{\mathrm{H}}\{a \in G : \ \dim_{\mathrm{H}}(J_A) < s(a)\} \leq \sup_G s(a).$$

(iii) J_A *has positive Lebesgue measure for for Lebesgue-a.e.* $a \in U_{>1}$.

Note that both $U_{<1}$ and $U_{>1}$ are non-empty. Indeed, continuity of $s(a)$ was established independent of transversality (see Lemma 3.4). Thus $s(a)$ is continuous on $(0, 2)$. For $a = \frac{1}{2}$ we know that the limit set of the IFS $\{\phi_1, \phi_3\}$ is an interval. Since adding an extra function to the IFS makes $s(a)$ strictly larger, we see that $s(\frac{1}{2}) > 1$. By continuity, this implies that $U_{>1} \neq \emptyset$. On the other hand, for $a = \sqrt{3}-1$ and for $a = 3-\sqrt{3}$ the IFS satisfies the open set condition, as the interiors of $\phi_i(Y)$, $i = 1, 2, 3$, are disjoint. Since the limit set is disconnected at these parameter values, it follows, as in [U, Th.6.5], that $s(\sqrt{3}-1) < 1$ and $s(3-\sqrt{3}) < 1$. By continuity, this implies that $U_{<1} \neq \emptyset$.

Example 4. Consider $\Phi_A = \{\phi(x + a_i)\}_{i \leq k}$ where

$$A \in U := \{A \in \mathbb{R}^k : \ a_k = 0, \ a_j > 0, \ j = 1, \ldots, k-1\}, \quad k \geq 3.$$

This is a $(k-1)$-parameter family of parabolic IFS.

Corollary 7.5. **(i)** *For Lebesgue-a.e.* $A \in U$,

$$\dim_{\mathrm{H}}(J_A) = \min\{s(A), 1\}.$$

(ii) *For any set $G \subset U_{<1}$ we have*

$$\dim_{\mathrm{H}}\{A \in G:\ \dim_{\mathrm{H}}(J_A) < s(A)\} \le \sup_G s(A) + (k-2).$$

(iii) *For Lebesgue-a.e. $A \in U_{>1}$ the set J_A has positive Lebesgue measure.*

Proof. As in Example 3, we have that if $a_i \le \frac{1}{2}$ for some $i \le k-1$, then J_A contains an interval an and there is nothing to prove. If $a_i > \frac{1}{2}$ for all i, then (7.7) holds. Then we can apply Proposition 7.2 (with $d = k-1$, when (7.2) is vacuous), and the statements follow from our theorems. □

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